# KÄHLER-EINSTEIN METRICS ON STABLE VARIETIES AND LOG CANONICAL PAIRS

by

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Abstract. — Let X be a canonically polarized variety, i.e. a complex projective variety such that its canonical class  $K_X$  defines an ample  $\mathbb{Q}$ —line bundle, and satisfying the conditions  $G_1$  and  $S_2$ . Our main result says that X admits a Kähler-Einstein metric iff X has semi-log canonical singularities i.e. iff X is a stable variety in the sense of Kollár-Shepherd-Barron and Alexeev (whose moduli spaces are known to be compact). By definition a Kähler-Einstein metric in this singular context simply means a Kähler-Einstein on the regular locus of X with volume equal to the algebraic volume of  $K_X$ , i.e. the top intersection number of  $K_X$ . We also show that such a metric is uniquely determined and extends to define a canonical positive current in  $c_1(K_X)$ . Combined with recent results of Odaka our main result shows that X admits a Kähler-Einstein metric iff X is K-stable, which thus confirms the Yau-Tian-Donaldson conjecture in this general setting of (possibly singular) canonically polarized varieties. More generally, our results are shown to hold in the setting of log minimal varieties and they also generalize some prior results concerning Kähler-Einstein metrics on quasi-projective varieties.

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## Introduction

According to the seminal works of Aubin [Aub78] and Yau [Yau78b] any canonically polarized compact complex manifold X (i.e. X is a non-singular projective algebraic variety such that the canonical line bundle  $K_X$  is ample) admits a unique Kähler-Einstein metric  $\omega$  in the first Chern class  $c_1(K_X)$ . One of the main goals of the present paper is to extend this result to the case when X is singular or more precisely when X has semi-log canonical singularities. A major motivation comes from the fact that such singular varieties appear naturally in the compactification of the moduli space of canonically polarized manifolds a subject where there has been great progress in the last years in connection to the (log) Minimal Model Program (MMP) in birational algebraic geometry [Kol, Kov12]. The varieties in question are usually referred to as stable varieties (or canonical models) as they are the higher dimensional generalization of the classical notion of stable curves of genus q > 1, which form the Deligne-Mumford compactification of the moduli space of non-singular genus q curves [Kol, Kov12]. It as a classical fact that any stable curve admits a unique Kähler-Einstein metric on its regular part, whose total area is equal to the (arithmetic) degree of the curve X and our first main result gives a generalization of this fact to the higher dimensional setting:

**Theorem A.** — Let X be a projective complex algebraic variety with semi-log canonical singularities such that  $K_X$  is ample. Then there exists a Kähler metric on the regular locus  $X_{reg}$ , satisfying

and such that the volume of  $(X_{reg}, \omega)$  coincides with the volume of  $K_X$ , i.e.  $\int_{X_{reg}} \omega^n = c_1(K_X)^n$ . Moreover, the metric extends to define a current  $\omega$  in  $c_1(K_X)$  which is uniquely determined by X.

We will refer to the current  $\omega$  in the previous theorem as a (singular) Kähler-Einstein metric on X. Moreover, the current  $\omega$  will be shown to be of *finite energy*, in the sense of [GZ07, BBGZ09] and as discussed in the last section of the present paper this allows one to define a canonical (singular) Weil-Peterson metric on the compact moduli space in terms of Deligne pairings. The notion of semi-log canonical singularities of a variety X - which is the most general class of singularities appearing in the (log) Minimal Model Program - will be recalled below. For the moment let us just point out that the definition involves two ingredients: first a condition which makes sure that the canonical divisor  $K_X$  is defined as a  $\mathbb{Q}$ - Cartier divisor (i.e.  $\mathbb{Q}$ -line bundle) which is in particular needed to make so sense of the notion of ampleness of  $K_X$  and secondly, the definition of semi-log canonical singularities involves a bound on the discrepancies of X on any resolution of singularities.

In fact, we will conversely show that if  $K_X$  is ample and the variety X admits a Kähler-Einstein metric then X has semi-log canonical singularities and this brings us to our second motivation for studying Kähler-Einstein metrics in the the singular setting, namely the Yau-Tian-Donaldson conjecture. Recall that this conjecture concerns polarized algebraic manifolds (X, L), i.e. algebraic manifolds together with an ample line bundle  $L \to X$  and it says that the first Chern class  $c_1(L)$  of L contains a Kähler metric  $\omega$  with constant scalar curvature if and only if (X, L) is K-stable. The latter notion of stability is of an algebrogeometric nature and can be seen as an asymptotic form of the classical notions of Chow and Hilbert stability appearing in Geometric Invariant Theory (GIT). However, while the notion of K-stability makes equal sense when X is singular it is less clear how to give a proper definition of a constant scalar curvature metric for a singular polarized variety (X, L). But, as it turns out, the situation becomes more transparent in the case when L is equal to  $K_X$  or its dual, the anti-canonical bundle  $-K_X$ . The starting point is the basic fact that, when X is smooth, a Kähler metric in  $\omega$  in  $c_1(\pm K_X)$  has constant scalar curvature on all of X precisely when it has constant Ricci curvature, i.e. when  $\omega$  is a Kähler-Einstein metric. Various generalizations of Kähler-Einstein metrics to the singular setting have been proposed in the litterature, see e.g. [EGZ09, BEGZ10, BBE+11] etc. In this paper we will adopt the definition which appears in the formulation of the previous theorem (see section 1), i.e. positive current in  $c_1(\pm K_X)$  is said to define a (singular) Kähler-Einstein metric if defines a bona fide Kähler-Einstein metric on the regular locus  $X_{\text{reg}}$  and if its total volume there coincides with the algebraic top intersection number of  $c_1(\pm K_X)$ . This definition, first used in the Fano case in [BBE+11], has the virtue of generalizing all previously proposed definitions, regardless of the sign of the canonical line bundle. Combing our results with recent results of Odaka [Oda08, Oda11], which say that a canonically polarized variety has semi-log canonical singularities precisely when  $(X, K_X)$  is K-stable, gives the following theorem, which can be seen as a confirmation of the generalized form of the Yau-Tian-Donaldson conjecture for canonically polarized varieties (satisfying the conditions  $G_1$  and  $S_2$ , cf [?]):

**Theorem B.** — Let X be a projective complex algebraic variety such that  $K_X$  is ample. Then X admits a Kähler-Einstein metric if and only if  $(X, K_X)$  is K-stable.

It may also be illuminating to compare this result with the case when  $L := -K_X$  is ample (i.e. X is Fano). Then it was shown in [Ber12], in the general singular setting, that the existence of a Kähler-Einstein metrics indeed implies K-(poly)stability. As for the converse it was finally settled very recently in the deep works by Chen-Donaldson-Sun [CDS12a, CDS12b, CDS13] and Tian [Tia13], independently, in the case when X is smooth. The existence problem in the singular case is still open in general, except for the toric case [BB12]; cf also [OSS12] for a related problem in the case of singular Fano surfaces.

Coming back to the present setting we point out that the starting point of our approach is that, after passing to a suitable resolution of singularities, we may as well assume that the variety X is smooth if we work in the setting of  $log\ pairs\ (X,D)$ , where D is a  $\mathbb{Q}$ -divisor on X with simple normal crossings (SNC) and where the role of the canonical line bundle is played by the  $log\ canonical\ line\ bundle\ K_X + D$  (which appears as the pull-back to the resolution of the original canonical line bundle). In this notation the original variety has semi-log canonical singularities precisely when the log pair (X,D) is  $log\ canonical\ (lc)$  in the usual sense of the Minimal Model Program, i.e. the coefficients of D are at most equal to one (but negative coefficients are allowed). However, it should be stressed that for this gain in regularity we have, of course, to pay a loss of positivity: even if the original canonical line bundle is ample, the corresponding log canonical line bundle is only semi-ample (and big) on the resolution, since it is trivial along the exceptional divisors of the corresponding resolution.

The upshot is that the natural setting for our results is the setting of log canonical pairs (X, D) such that the log canonical line bundle  $K_X + D$  is semi-ample and big. To any such pair we will associate a canonical Khler-Einstein metric  $\omega$  in the sense that  $\omega$  is a current in the first Chern class  $c_1(K_X + D)$  such that  $\omega$  restricts to a bona fide Kähler-Einstein metric on a Zariski open set of X and such that, globally on X, the current defined by the divisor D gives a singular contribution to the Ricci curvature of  $\omega$ .

The existence proof of Theorem A (and its generalizations described below) will be divided into two parts: in the first part we construct a variational solution with finite energy, by adapting the variational techniques developed in [BBGZ09] to the present setting. Then, in the second part, we show that the variational solutions have appropriate regularity using a priori Laplacian estimates, building on the works of Aubin [Aub78] and Yau [Yau78b] and ramifications of their work by Kobayashi [Kob84] and Tian-Yau [TY87] to the setting of quasi-projective varieties - in particular we will be relying on Yau's maximum principle. For

the second part we will need to perturb the line bundle  $L := K_X + D$  (to make it ample) and regularize the klt part of the divisor D (to make the divisor purely log canonical). To this end it will be convenient to consider the even more general setting of twisted Kähler-Einstein metrics attached to a twisted log canonical pairs  $(X, D, \theta)$  where (X, D) is a log canonical pair and  $\theta$  is a smooth form on X such that the twisted log canonical line bundle  $K_X + D + [\theta]$  is semi-positive and big (in the perturbation argument it will be ample/Kähler). Then the twisted Kähler-Einstein metric is defined as before, after replacing  $\text{Ric }\omega$  with the twisted Ricci curvature  $\text{Ric }\omega - \theta$ .

**Theorem C.** — Let X be a Kähler manifold and D a simple normal crossings  $\mathbb{R}$ —divisor on X with coefficients in  $]-\infty,1]$  such that  $K_X+D$  is semi-positive and big (i.e.  $(K_X+D)^n>0$ ). Then there exists a unique current  $\omega$  in  $c_1(K_X+D)$  which is smooth on a Zariski open set U of X and such that

$$\operatorname{Ric} \omega = -\omega + [D]$$

holds on X in the weak sense and  $\int_U \omega^n = (K_X + D)^n$ . More precisely, U can be taken to be the complement of D in the ample locus of  $K_X + D$ . Moreover,

- Any such current  $\omega$  on X automatically has finite energy.
- $\omega$  can be approximated by complete approximate Kähler-Einstein metrics  $\omega_j$  in the following sense: split  $D = D_{lc} + D_{klt}$  where  $D_{lc}$  is the the purely log canonical part and let  $\theta_j$  be a sequence of smooth forms regularizing the current  $D_{klt}$  and fix a Kähler form  $\omega_0$  on X. Then the twisted Kähler-Einstein metrics  $\omega_j$  attached to  $(X, D_{lc}, \theta_j + \frac{1}{j}\omega_0)$  are smooth and complete on  $X D_{lc}$  and  $\omega_j \to \omega$  in the weak toplogy of currents on X and in the  $C_{loc}^{\infty}$  topology on the Zariski open set U above.

The previous theorem will, more generally, be shown to hold (with essentially the same proof) as long as  $K_X + D$  is nef and big, see Theorem 4.11. This is indeed more general than assuming only semi-positivity unless D is klt, as the Shokurov-Kawamata base point free theorem does not hold for general lc pairs: there exists indeed a counterexample due to Zariski, cf [KMM87, Remark 3-1-2]. Let us also point out that the variational part of our proof only requires  $K_X + D$  to be big and thus produces a unique singular Kähler-Einstein metric of finite energy on any variety of log general type. As for the regularity part it applies as long as the corresponding log canonical ring is finitely generated. In fact, according to one of the fundamental conjectures of the general log minimal model program the latter finiteness property always holds. Note that for log canonical pairs, the finite generation is known to hold for  $n \leq 4$ , cf [K+92, Fuj10] and the references therein. We also recall that in the case of varieties of log general type with log terminal singularities (which in our notation means that  $D_{lc}$  vanishes) the finite generation in question was established in the seminal work [BCHM10] and Kähler-Einstein metrics were first constructed in [EGZ09] (cf also the discussion following Corollary 4.12).

In the last section of the paper some applications of Theorem A are given. First, we explain the link with Yau-Tian-Donaldson as we indicated above in Theorem B. Then, we

give a short analytic proof of the fact that the automorphism group of a canonically polarized variety with log canonical singularities is finite (see [BHPS12] for algebro-geometric proofs). We also discuss the problem of deducing Miyaoka-Yau type inequalities from Theorem A.

Further comparison with previous results. — The previous theorem also extends some of the results of Wu in [Wu08, Wu09], concerning the setting of Kähler-Einstein metrics on quasi-projective projective varieties of the form  $X_0 := X - D$ , where X is smooth and D is reduced SNC divisor. We recall that the case when  $K_X + D$  is ample was independently settled by Kobayashi [Kob84] and Tian-Yau [TY87]. The case when X is an orbifold and  $K_X + D$  is semi-ample and big was considered by Tian-Yau in [TY87] and as later shown by Yau [Yau93] the corresponding Kähler-Einstein metric is then complete on  $X_0$ . (in the orbifold sense). However, in our general setting the metric will typically not be complete on the regular locus. This is only partly due to the klt singularities (which generalize orbifold singularities) - there is also a complication coming from the presence of negative coefficients on a resolution (see section 4.8).

To illustrate this we recall that a standard example of log canonical pairs (X, D) is given by the Borel-Baily compactification  $X := X_0 - D$  of an arithmetic quotient, i.e.  $X_0 = B/\Gamma$ , where B is a bounded symmetric domain and  $\Gamma$  is discrete subgroup of the automorphism group of B. In this case any toroidal resolution X' has the property that the corresponding divisor D' on the resolution is reduced (and hence purely log canonical) if  $\Gamma$  is neat, i.e. if there are no fixed points. The corresponding Kähler-Einstein metric on  $X_0$  is the complete one induced from the corresponding metric on B, constructed in [CY80, MY83]. When  $\Gamma$  has fixed points these give rise to an additional fractional klt part  $D'_{klt}$  in D' so that the corresponding Kähler-Einstein metric is only complete in the orbifold sense [TY87]. However, for general log canonical singularity (X, D) the klt part  $D'_{klt}$  of D' may not be fractional or more seriously: it may contain negative coefficients and the main novelty of the present paper is to show how to deal with this problem by combining a variational approach with a priori estimates.

It is also interesting to compare with the case when the pair (X, D) is log smooth with  $K_X+D$  ample and with D effective and klt (i.e. with coefficients in [0,1[), where very precise regularity results have been obtained recently. For example, in [Bre11, CGP11, JMR11] it is shown that the corresponding Kähler-Einstein metric  $\omega$  has conical singularities along D (sometimes also called edge singularities in the litterature), thus confirming a previous conjecture of Tian – note that the results of [CGP11] holds for coefficients bigger than 1/2 whereas in [JMR11] the divisor D is assumed to be smooth but with no assumptions on the coefficients, and in [Bre11] both conditions are required; moreover the general case is announced in [MR12]. As for the mixed case when the coefficient 1 is also allowed in D it was studied in [Gue12b], where it was shown that  $\omega$  has mixed cone and Poincaré type singularities. A commun theme in these results is that singularities of the metric  $\omega$  are encoded by a suitable local model (with cone or Poincaré type singularities) determined by D. However, the difficulty in the situation studied in the present paper is the presence of

negative coefficients in D and the associated loss of positivity which appears when we pass to a log resolution of a singular variety X. It would be very interesting if one could associate local models to this situation as well, but this seems very challenging even in the case when X has klt singularities.

# Organization of the paper. —

- §1: We introduce the preliminary material that we will need, concerning the pluripotential theoretic setting of singular metrics on line bundles over varieties which are not necessarily normal.
- §2: Here we give the definition of a Kähler-Einstein metric on a canonically polarized variety X and more generally on a log pair (X, D). As we explain a purely differential-geometric definition can be given which only involves the regular locus  $X_{\text{reg}}$  of X. But, at we show, the corresponding metric automatically extends in a unique manner to define a singular current on X (which will allow us to prove the uniqueness of the Kähler-Einstein metric, later on in section 3). We first treat the case when X has log canonical (and hence normal singularities) and then the general case of a variety X with semi-log singularities. Anyway, as we recall, the latter case reduces to the former (if one works in the setting of pairs) if one passes to the normalization.
- §3: We prove the uniqueness and existence of a weak Kähler-Einstein metric in the general setting of varieties of log general type. The existence is proved by adapting the variational approach to complex Monge-Ampère equations introduced in [BBGZ09] to the present setting. This method produces a singular Kähler-Einstein metric with finite energy (the new feature here compared to [BBGZ09] is that the reference measure does not have an  $L^1$  density). We also use the variational approach to establish a stability result for the solutions to the equations induced from an (ample) perturbation of the log canonical line bundle on a resolution.
- §4: Here we establish the smoothness of the Kähler-Einstein metric, produced by the variational approach, on the regular locus of the variety X (or more generally, the pair (X, D)). The proof uses a perturbation argument in order to reduce the problem the the original setting of Kobayashi and Tian-Yau, combined with a priori estimates. But it should be stressed that in order to control the  $\mathcal{C}^0$  norms we need to invoke the variational stability result proved in the previous section. We also investigate the (non-) completeness properties of the Kähler-Einstein metrics.
- §5: We give some applications to automorphism groups and show how to deduce the Yau-Tian-Donaldson conjecture for canonically polarized varieties from our results.
- §6: The paper is concluded with a brief outlook on possible applications to Miyaoka-Yau types inequalities, as well as the Weil-Peterson geometry of the moduli space of stable varieties. These applications will require a more detailed regularity analysis of the Kähler-Einstein metrics that we leave for the future.

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## 1. Preliminaries

We collect here some useful tools or notions that we are going to work with in this paper. We start with a compact Kähler manifold X of dimension n, and we consider a class  $\alpha \in H^{1,1}(X,\mathbb{R})$  which is big. By definition, this means that  $\alpha$  lies in the interior of the pseudo-effective cone, so that there exists a Kähler current  $T \in \alpha$ , that is a current which dominates some smooth positive form  $\omega$  on X. We fix  $\theta$ , a smooth representative of  $\alpha$ .

The ample locus. — An important invariant attached to  $\alpha$  is the ample locus of  $\alpha$ , denoted  $Amp(\alpha)$ , and introduced in [Bou04, §3.5]. This is the largest Zariski open subset U of X such that for all  $x \in U$ , there exists a Kähler current  $T_x \in \alpha$  with analytic singularities such that  $T_x$  is smooth in an (analytic) neighbourhood of x. Its complement, called the augmented base locus, is usually denoted by  $\mathbb{B}_+(\alpha)$ . In the case when  $\alpha = c_1(L)$  is the Chern class of a line bundle, it is known (see e.g. [BBP10]) that:

$$\mathbb{B}_{+}(L) = \bigcap_{\substack{L=A+E\\A \text{ ample}, E \geqslant 0}} \operatorname{Supp}(E)$$

Currents with minimal singularities. — We will be very brief about this well-known notion, and refer e.g. to [Bou04, §2.8], [BBGZ09, §1], [Ber09] or [BD12] for more details and recent results.

By definition, if T, T' are two positive closed currents in the same cohomology class  $\alpha$ , we say that T is less singular than T' if the local potentials  $\varphi, \varphi'$  of these currents satisfy  $\varphi' \leqslant \varphi + O(1)$ . It is clear that this definition does not depend on the choice of the local potentials, so that the definition is consistant. In each (pseudo-effective) cohomology class  $\alpha$ , one can find a positive closed current  $T_{\min}$  which will be less singular than all the other ones; this current is not unique in general; only its class of singularities is. Such a current will be called current with minimal singularities.

One way to find such a current is to pick  $\theta \in \alpha$  a smooth representative, and define then, following Demailly, the upper envelope

$$V_{\theta} := \sup \{ \varphi \ \theta - \text{psh}, \ \varphi \leqslant 0 \text{ on } X \}$$

Once observed that  $V_{\theta}$  is  $\theta$ -psh (in particular upper semi-continuous), it becomes clear that  $\theta + dd^cV_{\theta}$  has minimal singularities.

Non-pluripolar Monge-Ampère operator. — In the paper [BEGZ10], the authors define the non-pluripolar product  $T \mapsto \langle T^n \rangle$  of any closed positive (1,1)-current  $T \in \alpha$ , which is shown to be a well-defined measure on X putting no mass on pluripolar sets, and extending the usual Monge-Ampère operator for Kähler forms (or having merely bounded potentials, cf [BT87]). Let us note that when T is a smooth positive form  $\omega$  on a Zariski dense open subset  $\Omega \subset X$ , then its Monge-Ampère  $\langle T^n \rangle$  is simply the extension by 0 of the measure  $\omega^n$  defined on  $\Omega$ .

Given now a  $\theta$ -psh function  $\varphi$ , one defines its non-pluripolar Monge-Ampère by  $\mathrm{MA}(\varphi) := \langle (\theta + dd^c \varphi)^n \rangle$ . Then one can check easily from the construction that the total mass of  $\mathrm{MA}(\varphi)$  is less than or equal to the volume  $\mathrm{vol}(\alpha)$  of the class  $\alpha$  (cf [Bou02]):

$$\int_X \mathrm{MA}(\varphi) \leqslant \mathrm{vol}(\alpha)$$

A particular class of  $\theta$ -psh functions that appears naturally is the one for which the last inequality is an equality. We will say that such functions (or the associated currents) have full Monge-Ampère mass. For example,  $\theta$ -psh functions with minimal singularities have full Monge-Ampère mass (cf [BEGZ10, Theorem 1.16]).

Plurisubharmonic functions on complex spaces. — Here again, we just intend to give a short overview of the extension of the pluripotential theory to (reduced) complex Kähler spaces. A very good reference is [Dem85], or [EGZ09, §5] which is written in relation to singular Kähler-Einstein metric. We also refer to the preliminary parts of [Var89] or [FS90].

The data of a reduced complex space X includes the data of the sheaves of continuous and holomorphic functions. So the first object we would like to give a sense to is the sheaf  $\mathscr{C}_X^{\infty}$  of smooth functions. It may be defined as the restriction of smooth functions in some local embeddings of X in some  $\mathbb{C}^n$ . One defines similarly the sheaves of smooth (p,q)-forms  $\mathscr{A}_X^{p,q}$  which carry the differentials  $d, \partial, \bar{\partial}$  satisfying the usual rules; the space of currents is by definition the dual of the space of differential forms as in the smooth case. The sheaves complexes that are induced (Dolbeault, de Rham, etc.) are however not exact in general.

Another important sheaf is the one of pluriharmonic functions. They are defined to be smooth functions locally equal to the imaginary part of some holomorphic functions. One can show (see e.g [FS90]) that a continuous function which is pluriharmonic on  $X_{\text{reg}}$  in the usual sense is automatically pluriharmonic on X. We denote by  $\mathcal{PH}_X$  the sheaf of real-valued pluriharmonic functions on X.

Let us move on to psh functions now. There are actually two possible definitions which extend the usual one for complex manifolds. The first one, introduced by Grauert and Remmert, mimics the one in the smooth case: we will say that a function  $\varphi: X \to \mathbb{R} \cup \{-\infty\}$  is plurisubharmonic if it is upper semi-continuous and if for all holomorphic map  $f: \Delta \to X$  from the unit disc in  $\mathbb{C}$ , the function  $\varphi \circ f$  is subharmonic.

We could also introduce a more local definition: a function  $\varphi: X \to \mathbb{R} \cup \{-\infty\}$  is strongly plurisubharmonic if in any local embeddings  $i_{\alpha}: X \supset U_{\alpha} \hookrightarrow \mathbb{C}^{n}$ ,  $\varphi$  is the restriction of a psh function defined an open set  $\Omega_{i} \subset \mathbb{C}^{n}$  containing  $i_{\alpha}(U_{\alpha})$ , if  $X = \bigcup_{\alpha} U_{\alpha}$  is an open covering.

Clearly, a strongly psh function is also psh. Actually, Fornaess and Narasimhan [FN80] showed that these notions coincide: a function on X is psh if and only if it is strongly psh. On *normal* spaces one still has a Riemann extension theorem for psh functions, thanks to [GR56]. More precisely, if X is normal,  $Y \subseteq X$  is any proper analytic subspace, and  $\varphi: X \setminus Y \to \mathbb{R} \cup \{-\infty\}$  is psh, then  $\varphi$  extends to a (unique) psh function on X if and only if it is locally bounded above near the points of Y, condition which is always realized if Y has codimension at least two in X. In particular, if X is normal, the data of a psh function on X is equivalent to the data of a psh function on  $X_{\text{reg}}$ .

Moreover, one can show (cf [BEG13, Lemma 3.6.1]) that a pluriharmonic function on  $X_{\text{reg}}$  automatically extends to a pluriharmonic function on X.

On non-normal spaces, one has to be more cautious, and it is convenient to introduce the notion of weakly psh function. Let X be a reduced complex space, and  $\nu: X^{\nu} \to X$  its normalization. We say that a function  $\varphi: X \to \mathbb{R} \cup \{-\infty\}$  is weakly psh if  $\nu^* \varphi = \varphi \circ \nu$  is psh. It is not hard to see that a weakly psh function  $\varphi$  induces a bona fide psh function on  $X_{\text{reg}}$  which is locally bounded from above near the points of  $X_{\text{sing}}$ . Conversely, any psh function on  $X_{\text{reg}}$  which is locally upper bounded extends to a weakly psh function on X. On a normal space, a weakly psh function is of course psh, but in general these notions are different: consider  $X = \{zw = 0\} \subset \mathbb{C}^2$ , and  $\varphi(x) = 0$  or 1 according to the connected component of  $x \in X$ . We refer to [Dem85, Théorème 1.10] for equivalent characterizations of weakly psh functions and conditions on a weakly psh function that ensure that it is already psh.

Finally, one can check that a (strongly) psh function  $\varphi$  on a complex space X is always locally integrable with respect to the area measure induced by any local embedding of X in  $\mathbb{C}^n$  (note that this is stronger than saying that  $\varphi$  is locally integrable on  $X_{\text{reg}}$  with respect to some volume form). Moreover, a locally integrable function  $\varphi$  is (almost everywhere) weakly psh is and only if it is locally bounded from above and  $dd^c\varphi$  is a positive current.

Weights and Chern classes. — From now on, X will be a *normal* complex space unless stated otherwise.

The definition of a (smooth) Kähler form is rather natural: it is a smooth real (1,1)-form written locally as  $dd^c\psi$  for some (smooth) strictly psh function  $\psi$ ; equivalently this is locally the restriction of a Kähler form in a embedding in  $\mathbb{C}^n$ . Note that we could interpret this definition in terms of hermitian metrics on the Zariski tangent bundle of X, cf [Var89].

Let us now consider a line bundle L on X. A smooth hermitian metric h on L is defined as in the smooth case: using trivialisations  $\tau_{\alpha}: L_{|U_{\alpha}} \xrightarrow{\simeq} U_{\alpha} \times \mathbb{C}$ , we just ask h to be written as  $h(v) = |\tau_{\alpha}(v)|^2 e^{-\varphi_{\alpha}(z)}$  where  $\varphi_{\alpha}$  is a smooth function on  $U_{\alpha}$ . We say that the

data  $\phi := \{(U_{\alpha}, \varphi_{\alpha})\}\$  is a weight on L, so that it is equivalent to consider a weight or an hermitian metric.

Observe that if  $(g_{\alpha\beta}: U_{\alpha} \cap U_{\beta} \to \mathbb{C}^*)$  is the cocycle in  $H^1(X, \mathcal{O}_X^*)$  determined by the  $\tau_{\alpha}$ 's (more precisely  $\tau_{\alpha} \circ \tau_{\beta}^{-1}(z, v) = (z, g_{\alpha\beta}v)$ ), then we have necessarily  $\varphi_{\beta} - \varphi_{\alpha} = \log|g_{\alpha\beta}|^2$ . In particular, the forms  $dd^c\varphi_{\alpha}$  glue to a global smooth (1, 1)-form on X called curvature of (L, h) and denoted by  $c_1(L, h)$ . This forms lives naturally in the space  $H^0(X, \mathscr{C}_X^{\infty}/\mathcal{PH}_X)$ , and using the exact sequence

$$0 \longrightarrow \mathcal{PH}_X \longrightarrow \mathscr{C}_X^{\infty} \longrightarrow C_X^{\infty}/\mathcal{PH}_X \longrightarrow 0$$

one may attach to (L, h) a class  $\hat{c}_1(L, h) \in H^1(X, \mathcal{PH}_X)$ . It is then easy to see that this class actually does not depend on the choice of h, so we will denote it by  $\hat{c}_1(L)$ . If X is smooth,  $H^1(X, \mathcal{PH}_X) \simeq H^{1,1}(X, \mathbb{R})$ , and it is well-known that  $\hat{c}_1(L)$  coincides with the image of  $L \in H^1(X, \mathcal{O}_X^*)$  in  $H^2(X, \mathbb{Z})$  via the connecting morphism induced by the exponential exact sequence

$$0 \longrightarrow \mathbb{Z} \longrightarrow \mathcal{O}_X \stackrel{e^{2i\pi \cdot}}{\longrightarrow} \mathcal{O}_X^* \longrightarrow 0$$

This sequence also exists on any (even non-reduced) complex space, so that  $c_1(L) \in H^2(X, \mathbb{Z})$  is well-defined; it will be more convenient for us to look at the image of  $c_1(L)$  in  $H^2(X, \mathbb{R})$  however. To relate it to  $\hat{c}_1(L)$ , we may use the following exact sequence:

$$(1.1) 0 \longrightarrow \mathbb{R} \longrightarrow \mathcal{O}_X \stackrel{-2\operatorname{Im}(\cdot)}{\longrightarrow} \mathcal{P}\mathcal{H}_X \longrightarrow 0$$

It is not hard to check that the connecting morphism  $H^1(X, \mathcal{PH}_X) \to H^2(X, \mathbb{R})$  sends  $\hat{c}_1(L)$  to  $c_1(L)$  as expected.

We will also have to consider singular weights, which are by definition couples  $\phi := \{(U_{\alpha}, \varphi_{\alpha})\}$  where  $U_{\alpha}$  is covering of X trivializing L, and  $\varphi_{\alpha}$  are locally integrable on  $U_{\alpha}$ , satisfying  $\varphi_{\beta} - \varphi_{\alpha} = \log |g_{\alpha\beta}|^2$  on  $U_{\alpha} \cap U_{\beta}$ . The associated curvature current, denoted by  $dd^c\phi$ , is well-defined on X. The weight is said psh if the  $\varphi_{\alpha}$  are, in which case  $dd^c\phi$  is a positive current. Moreover, we can proceed as in the smooth case to attach to  $dd^c\phi$  a class  $\hat{c}_1(L) \in H^1(X, \mathcal{PH}_X)$  (consider  $\phi$  as a section of the sheaf  $L^1_{\text{loc}}/\mathcal{PH}_X$  and use the natural exact sequence), whose image in  $H^2(X, \mathbb{R})$  via the long exact sequence in cohomology induced by (1.1) is  $c_1(L)$ . Therefore when  $\phi$  is a singular weight on L, we may say that  $dd^c\phi$  is a current in  $c_1(L)$ .

We claim that a (possibly singular) psh weight  $\phi$  on  $L_{|X_{\text{reg}}}$  - and thus a psh weight in the usual sense- automatically extends to a (unique) psh weight  $\widetilde{\phi}$  on L. Indeed, by Grauert and Remmert's theorem, the  $\varphi_{\alpha}$ 's defined on  $U_{\alpha} \cap X_{\text{reg}}$  extend to a psh function  $\widetilde{\varphi_{\alpha}}$  on the whole  $U_{\alpha}$ , which is moreover defined by  $\widetilde{\varphi_{\alpha}}(z_0) = \limsup_{X_{\text{reg}} \ni z \to z_0} \varphi_{\alpha}(z)$ . Therefore, the relation  $\widetilde{\varphi_{\beta}} = \widetilde{\varphi_{\alpha}} + \log |g_{\alpha\beta}|^2$  is immediately satisfied on the whole  $U_{\alpha} \cap U_{\beta}$ , which proves the claim.

If we get back to the non-normal case, we can define psh weights using weakly psh functions instead of psh functions. So the general philosophy is that we always pull-back our objects to the normalization where things behave better, and the notions downstairs are defined and studied upstairs. For example, we can define on a normal variety the analogue of the non-pluripolar product and consider Monge-Ampère equations as in the smooth case (cf [BBE+11, §1.1-1.2]. Then, if we write them on a non-normal variety, they have to be thought as pulled-back to the normalization.

**Log canonical pairs.** — Following the by now common terminology of Mori theory and the minimal model program (cf e.g. [KM98]), a pair (X, D) is by definition a complex normal projective variety X carrying a Weil  $\mathbb{Q}$ -divisor D (not necessarily effective). We will say that the pair (X, D) is a log canonical pair if  $K_X + D$  (which is a priori defined as a Weil divisor) is  $\mathbb{Q}$ -Cartier, and if for some (or equivalently any) log resolution  $\pi: X' \to X$ , we have:

$$K_{X'} = \pi^*(K_X + D) + \sum a_i E_i$$

where  $E_i$  are either exceptional divisors or components of the strict transform of D, and the coefficients  $a_i$  satisfy the inequality  $a_i \ge -1$ .

#### 2. Singular Kähler-Einstein metrics

**2.1.** Kähler-Einstein metrics on pairs. — In this section, we will consider log pairs (X, D) where X is a complex normal projective variety, D is a Weil divisor, and  $K_X + D$  is assumed to be  $\mathbb{Q}$ -Cartier. We choose a psh weight  $\phi_D$  on  $X_{\text{reg}}$  satisfying  $dd^c\phi_D = [D_{|X_{\text{reg}}}]$ . The first definition concerns the Ricci curvature of currents:

**Definition 2.1.** — Let  $\omega$  be a positive current on  $X_{\text{reg}}$ ; we say that  $\omega$  is admissible if it satisfies:

- 1. Its non-pluripolar product  $\langle \omega^n \rangle$  defines a (locally) absolutely continuous measure on  $X_{\text{reg}}$  with respect to  $d\mathbf{z} \wedge d\bar{\mathbf{z}}$ , where  $\mathbf{z} = (z_i)$  are local holomorphic coordinates.
- 2. The function  $\log(\langle \omega^n \rangle / d\mathbf{z} \wedge d\bar{\mathbf{z}})$  belongs to  $L^1_{\text{loc}}(X_{\text{reg}})$ .

In that case, we define (on  $X_{\text{reg}}$ ) the Ricci curvature of  $\omega$  by setting  $\text{Ric}\,\omega := -dd^c\log\langle\omega^n\rangle$ .

Another way of thinking of this is to interpret the positive measure  $\langle \omega^n \rangle_{|X_{\text{reg}}}$  as a singular metric on  $-K_{X_{\text{reg}}}$  whose curvature is  $\text{Ric}\,\omega$  by definition.

The measure  $e^{\phi}$  for  $\phi$  a weight on  $K_X$ . — In the same spirit, we will use the convenient but somehow abusive notation  $e^{\phi}$  for  $\phi$  a weight on  $K_X$  (whenever the latter is defined as a  $\mathbb{Q}$ -line bundle) to refer to the positive measure  $e^{\varphi_{\mathbf{z}}}d\mathbf{z} \wedge d\bar{\mathbf{z}}$  defined on  $X_{\text{reg}}$  and extended by 0 to X; where  $\varphi_{\mathbf{z}}$  is the expression on some trivializing chart of  $X_{\text{reg}}$  (and hence of  $K_{X_{\text{reg}}}$  too) of  $\phi$ . In particular, for  $\phi$  a psh weight on  $K_X + D$ , the measure  $e^{\phi - \phi_D}$  can easily be pulled-back to any log resolution  $(\pi, X', D')$  of (X, D) (we pull it back over  $X_{\text{reg}} \setminus \text{Supp}(D)$ 

and then extend it by 0), where it become  $e^{\phi \circ \pi - \phi_{D'}}$ .

We may now introduce the notion of (negatively curved) Kähler-Einstein metric attached to a pair (X, D):

**Definition 2.2.** — Let (X, D) be a log pair; we say that a positive admissible current  $\omega$  is a Kähler-Einstein metric with negative curvature for (X, D) if:

- 1. Ric  $\omega = -\omega + [D]$  on  $X_{\text{reg}}$ ,
- 2.  $\int_{X_{\text{reg}}} \langle \omega^n \rangle = c_1 (K_X + D)^n$ .

This conditions are sufficient to show that a Kähler-Einstein metric is a global solution of a Monge-Ampère equation. More precisely, we have the following:

**Proposition 2.3.** Let (X, D) be a log pair, and  $\omega$  be a Kähler-Einstein metric for (X, D). Then  $\phi := \log \langle \omega^n \rangle + \phi_D$  extends to X as a psh weight with full Monge-Ampère mass on  $K_X + D$ , solution of

$$\langle (dd^c\phi)^n \rangle = e^{\phi - \phi_D}$$

Conversely, any psh weight  $\phi$  on  $K_X + D$  with full Monge-Ampère mass solution of the equation induces a Kähler-Einstein metric  $\omega := dd^c \phi$  for (X, D).

*Proof.* — On  $X_{\text{reg}}$ , we have  $dd^c\phi = \omega$  thus  $\phi$  is a psh weight on  $(K_X + D)_{|X_{\text{reg}}}$ , and thanks to a theorem of Grauert and Remmert, it extends through  $X_{\text{sing}}$  which has codimension at least 2. Clearly, we have  $\omega = dd^c\phi$  on X, and by condition 3. in the definition of a Kähler-Einstein metric,  $\phi$  has full Monge-Ampère mass. Then by definition, the two (non-pluripolar) measures  $\langle (dd^c\phi)^n \rangle$  and  $e^{\phi-\phi_D}$  coincide.

For the converse, let  $\omega := dd^c \phi$ ; clearly 1. and 3. are satisfied. Moreover,  $\phi$  and  $\phi_D$  are locally integrable, so that  $\omega$  is admissible and  $\text{Ric }\omega = -dd^c(\phi - \phi_D) = -\omega + [D]$ .

This proposition shows that the different definitions of what should be a singular Kähler-Einstein metric, appearing e.g. in [Ber11, BEGZ10, CGP11, EGZ09] etc. coincide. Moreover, one could equally define positively curved Kähler-Einstein metrics in an equivalent way as in [BBE+11]. In particular this objects, intrinsically defined on X, can also be seen on any log resolution in the usual way; in practice, we will most of the time work on log resolutions when dealing with existence or smoothness questions.

Note also that we could have chosen to define a Kähler-Einstein metric attached to a pair (X, D) (say satisfying  $K_X + D$  ample) to be a *smooth* Kähler metric  $\omega$  on  $X_{\text{reg}} \setminus \text{Supp}(D)$  which extends to an admissible current on  $X_{\text{reg}}$  satisfying there  $\text{Ric }\omega = -\omega + [D]$  and the mass condition  $\left(\int_{X_{\text{reg}}} \langle \omega^n \rangle = \right) \int_{X_{\text{reg}} \setminus \text{Supp}(D)} \omega^n = c_1(K_X + D)^n$ .

Then, our regularity Theorem (say combined with Proposition 5.1) shows a posteriori that this definition would have coincided with Definition 2.2.

Let us also mention that in the case of a log smooth log canonical pair (X, D), the same proof as [Gue12b, Proposition 2.5] shows that the data of a negatively curved Kähler-Einstein on (X, D) is equivalent to giving an admissible current  $\omega$  on  $X \setminus \text{Supp}(D)$  such that:

- · Ric  $\omega = -\omega$  on  $X \setminus \text{Supp}(D)$ ,
- · There exists C > 0 such that

$$C^{-1}dV \leqslant \prod_{a_j < 1} |s_j|^{2a_j} \cdot \prod_{a_k = 1} (|s_k|^2 \log^2 |s_k|^2) \ \omega^n \leqslant CdV$$

for some volume form dV on X, and where  $D = \sum a_i D_i$ ,  $s_i$  being a defining section of  $D_i$ .

**2.2.** Kähler-Einstein metrics on stable varieties. — Stable varieties, as considered e.g. in [KSB88, Kar00, Kol, Kov12] are the appropriate singular varieties to look at if one wants to compactify the moduli space of canonically polarized projective varieties (cf also [Vie95]). Before giving the precise definition of a stable variety, we explain very briefly that notion and give the connection with Kähler-Einstein theory. In the next section, we will give a more detail account of the type of singularities involved.

So first of all, we will consider complex varieties that are Gorenstein in codimension 1 (this condition replaces regularity in codimension 1 for normal varieties) and satisfy the condition  $S_2$  of Serre. Basically, the singularities in codimension 1 of our varieties are those of the union of two coordinate hyperplanes ("double crossing"). Now we want to recast them in the context given by the singularities of the minimal model program (MMP); so we consider such a variety X and its normalization  $\nu: X^{\nu} \to X$ . One can write  $\nu^* K_X = K_{X^{\nu}} + D$  for some reduced divisor D called the conductor of  $\nu$ ; its sits above the codimension 1 component of the singular locus of X. We then say that X has semi log canonical singularities if the pair  $(X^{\nu}, D)$  is log canonical in the usual sense. The generalization of the notion of stable curve is given by the following definition:

**Definition 2.4.** — A projective variety X is called stable if X has semi-log canonical singularities, and  $K_X$  is an ample  $\mathbb{Q}$ -line bundle.

There is a subtlety for the definition of  $K_X$ , but we refer to §2.3 for appropriate explanations. Its is actually possible to define the notion of Kähler-Einstein metric for a stable variety:

**Definition 2.5.** — Let X be a stable variety. A Kähler-Einstein metric on X is a positive admissible current  $\omega$  on  $X_{\text{reg}}$  such that:

- 1. Ric  $\omega = -\omega$  on  $X_{\text{reg}}$ ,
- 2.  $\int_{X_{\text{reg}}} \langle \omega^n \rangle = c_1(K_X)^n$ .

In the non-normal case however, psh weight do not automatically extend across the singularities, so that it is not clear that the Kähler-Einstein metric will extend as a positive current on  $K_X$  satisfying the usual Monge-Ampère equation globally. Actually, this is the case as shows the following proposition:

**Proposition 2.6.** Let X be a stable variety, and  $\omega$  a Kähler-Einstein metric on X. Then the weight  $\phi := \log \omega^n$  extends to X as a weakly psh weight on  $K_X$  solution of the Monge-Ampère equation  $\langle (dd^c \phi)^n \rangle = e^{\phi}$ .

Proof. — Taking the  $dd^c$  of each side in the definition of  $\phi$  and using the Ricci equation, we find  $\omega = dd^c\phi$  on  $X_{\rm reg}$ , and therefore  $\phi$  satisfies  $\langle (dd^c\phi)^n \rangle = \langle \omega^n \rangle = e^{\phi}$ . Pulling back this equation to normalization  $X^{\nu}$ , we find a psh weight  $\phi' = \nu^*\phi$  on  $c_1(\nu^*K_X)_{|\nu^{-1}(X_{\rm reg})}$  solution of  $\langle (dd^c\phi')^n \rangle = e^{\phi'-\phi_D}$  where D is the conductor of the normalization. As we work inside  $X^{\nu}_{reg}$  and the integral  $\int_{\nu^{-1}(X_{\rm reg})} e^{\phi'-\phi_D}$  is finite, we infer from Lemma 2.7 below that  $\phi'$  extends (as a psh weight) across  $D_{\rm reg}$ . So  $\phi'$  induces a psh weight on  $c_1(\nu^*K_X)_{|X^{\nu}_{\rm reg}\setminus D_{\rm sing}}$ , and by normality of  $X^{\nu}$ , it extends to the whole  $X^{\nu}$ , which means precisely that  $\phi$  extends as a weakly psh weight on  $K_X$ . The expected Monge-Ampère equation holds automatically on X (or equivalently on X') since both measures  $\langle (dd^c\phi)^n \rangle$  and  $e^{\phi}$  put no mass on  $X_{\rm sing}$  by definition.

In the previous proof, we used the following extension result:

**Lemma 2.7.** Let U be a neighbourhood of  $0 \in \mathbb{C}^n$ ,  $H = \{z_1 = 0\} \subset \mathbb{C}^n$ , and  $\varphi$  be a psh function defined on  $U \setminus H$ . We assume that the integral

$$\int_{U\backslash H} \frac{e^{\varphi}}{|z_1|^2} dV$$

is finite. Then  $\varphi$  extends across H, and more precisely  $\varphi$  tends to  $-\infty$  near H.

Proof. — (thanks to Bo Berndtsson for providing us with this elegant proof) Assume, to get a contradiction, that  $\varphi$  does not tend to  $-\infty$  near H, and let  $V := U \setminus H$ . Then we can find a sequence  $(x_k)$  of points in V converging to H such that  $\varphi(x_k) \geqslant -C$  for some constant C. Moreover, one may assume that  $x_k = (t_1^k, \ldots, t_n^k)$  converges to 0, and that  $|t_1^{k+1}| < |t_1^k|/3$ . Thus we can find some polydiscs  $D_k$  centered at  $x_k$ , with polyradius  $(r_k, \delta, \ldots, \delta)$  such that  $r_k = |t_1^k|/2$ , and  $\delta > 0$  is fixed; in particular a simple computations shows that no two of the  $D_k$  meet. Now, using the mean value inequality for  $\varphi$  at  $x_k$ , we find:

$$-C \leqslant \frac{1}{\operatorname{vol}(D_k)} \int_{D_k} \varphi \, dV$$

Therefore, using Jensen's inequality, we obtain, up to modifying C by a normalization factor depending only on the dimension n:

$$e^{-C} \leqslant \int_{D_k} \frac{e^{\varphi} dV}{r_k^2 \delta^{2(n-1)}}$$

but on  $D_k$ ,  $|z_1| \leq 3r_k$  so

$$e^{-C'} \leqslant \int_{D_r} \frac{e^{\varphi} dV}{|z_1|^2}$$

for  $C' = C + \log 9 - 2(n-1) \log \delta$ . As the  $D_k$ 's are disjoint, it shows that the integral  $\int_{U \setminus H} \frac{e^{\varphi}}{|z_1|^2} dV$  is infinite, which is absurd.

**2.3.** Singularities of stable varieties. — In this paragraph, we intend to give a more precise overview of the notion of semi-log canonical singularities. As we will just touch on this topic, we refer to the nice survey [Kov12] for a broader study. Other good references are [Kol, KSB88].

In the following, X will always be a reduced and equidimensional scheme of finite type over  $\mathbb{C}$ , and we set  $n := \dim X$ .

The conditions  $G_1$  and  $S_2$ . — As we saw earlier, we need a canonical sheaf. The condition  $G_1$  will guarantee its existence, and the condition  $S_2$  will (among other things) ensure its uniqueness.

If X is Cohen-Macaulay (for every  $x \in X$ , the depth of  $\mathcal{O}_{X,x}$  is equal to its Krull dimension), then X admits a dualizing sheaf  $\omega_X$ . We say that X is Gorenstein if X is Cohen-Macaulay and  $\omega_X$  is a line bundle. We say that X is  $G_1$  if X is Gorenstein in codimension 1, which means that there is an open subset  $U \subset X$  which is Gorenstein and satisfies  $\operatorname{codim}_X(X \setminus U) \geq 2$ .

We say that X satisfies the condition  $S_2$  of Serre if for all  $x \in X$ , we have  $\operatorname{depth}(\mathcal{O}_{X,x}) \ge \min\{\operatorname{ht}(\mathfrak{m}_{X,x}), 2\}$ , where  $\operatorname{ht}(\mathfrak{m}_{X,x}) = \operatorname{codim}(\bar{x})$  denotes the height of the maximal ideal  $\mathfrak{m}_{X,x}$  of  $\mathcal{O}_{X,x}$ . This condition is equivalent to saying that for each closed subset  $i: Z \hookrightarrow X$  of codimension at least two, the natural map  $\mathcal{O}_X \to i_* \mathcal{O}_{X \setminus Z}$  is an isomorphism.

If X is  $G_1$  and  $S_2$ , and  $U \subset X$  is a Gorenstein open subset whose complement has codimension at least 2, one can then define the canonical sheaf  $\omega_X$  by  $\omega_X := j_*\omega_U$  where  $j: U \hookrightarrow X$  is the open embedding, and  $\omega_U$  is the dualizing sheaf of U. By definition, this is a rank one reflexive sheaf. When X is projective, we know that it admits a dualizing sheaf; as it is reflexive, it coincides with  $\omega_X$  by the  $S_2$  condition.

We would like to have an interpretation of  $\omega_X$ , or at least  $\omega_U$  in terms of Weil divisor as in the normal case where we define the Weil divisor  $K_X$  as the closure of some Weil divisor representing the line bundle  $K_{X_{\text{reg}}}$ . But we have to be more cautious in the non normal case it is not clear how we should extend a Weil divisor given on  $X_{\text{reg}}$ . Actually, this is where the  $G_1$  conditions appears: as  $\omega_U$  is a line bundle, or equivalently a Cartier divisor, we may choose a Weil divisor  $K_U$  whose support does not contain any component of  $X_{\text{sing}}$  of codimension 1 and represent  $\omega_U$  (write  $\omega_U$  as the difference of two very ample line bundles). Then we define  $K_X$  to be the closure of  $K_U$ . Clearly, the divisorial sheaf  $\mathcal{O}_X(K_X)$  is reflexive, and coincides with  $\omega_U = \omega_{X|U}$  on U, so that by the  $S_2$  condition, we get:

$$\omega_X \simeq \mathcal{O}_X(K_X)$$

In fact, if  $\omega_X^{[m]}$  denotes the m-th reflexive power of  $\omega_X$ , the same arguments yield  $\omega_X^{[m]} \simeq \mathcal{O}_X(mK_X)$ . Therefore, the Weil divisor is  $\mathbb{Q}$ -Cartier if and only if  $\omega_X$  is a  $\mathbb{Q}$ -line bundle, ie  $\omega_X^{[m]}$  is a line bundle for some m > 0.

Conductors and slc singularities. — Let now X be a (reduced) scheme, and  $\nu: X^{\nu} \to X$  its normalization. We recall that if X is not irreducible, its normalization is defined to be the disjoint union of the normalization of its irreducible components. The conductor ideal

$$\operatorname{cond}_X := \mathscr{H}om_{\mathcal{O}_X}(\nu_*\mathcal{O}_{X^{\nu}}, \mathcal{O}_X)$$

is the largest ideal sheaf on X that is also an ideal sheaf on  $X^{\nu}$ . If we think of the case where B is the integral closure of some integral ring A, then we can easily see that  $\operatorname{Hom}_A(B,A)$  injects in A (via the evaluation at 1), and the image of this map is the annihilator  $\operatorname{Ann}_A(B/A) = \{f \in A; fB \subset A\}$ , or equivalently the largest ideal  $\mathcal{I} \subset A$  that is also an ideal in B.

Coming back to the case of varieties, we will denote by  $\operatorname{cond}_{X^{\nu}}$  the conductor seen as an ideal sheaf on  $X^{\nu}$ , and we define the conductor subschemes as  $C_X := \operatorname{Spec}_X(\mathcal{O}_X/\operatorname{cond}_X)$  and  $C_{X^{\nu}} := \operatorname{Spec}_{X^{\nu}}(\mathcal{O}_{X^{\nu}}/\operatorname{cond}_{X^{\nu}})$ . If X is  $S_2$ , then one can show that these schemes have pure codimension 1 (and hence define Weil divisors) but they are in general not reduced (e.g. the cusp  $y^2 = x^3$ ).

If  $K_X$  is  $\mathbb{Q}$ -Cartier and X is demi-normal (ie X is  $S_2$  and has only double crossing singularities in codimension 1, cf [Kol]), we have the following relation:

$$(2.1) \nu^* K_X = K_{X^{\nu}} + C_{X^{\nu}}$$

The proof of this identity goes as follows: first, using the demi-normality assumption, we may assume that the only singularities of X are double normal crossings. Then, using the universal property of the dualizing sheaf (which coincide with the canonical sheaf as we observed above) and the projection formula, we have  $\nu_*\omega_{X^{\nu}}=\omega_X(-C_X)$ . We pull-back this relation to  $X^{\nu}$  using the fact that the sheaf  $\mathcal{O}_X(-C_X)$  becomes precisely  $\mathcal{O}_{X^{\nu}}(-C_{X^{\nu}})$ . By the assumptions on the singularities, this last sheaf is actually an invertible sheaf so that we get the expected identity (cf point 8 in [Kol]). As we will explain below, we do not want to assume a priori that our varieties are demi-normal. Therefore, it may happen that  $C_{X^{\nu}}$  is not Cartier, and the formula (2.1) may not be true anymore. So whenever we will deal with Kähler-Einstein on those varieties, we will have to apply the arguments on a log-resolution of the normalization instead of the normalization itself. Anyway, this will not cause any troubles.

An important point is that whenever the conductor is reduced, then necessarily X is seminormal (ie every finite morphism  $X' \to X$  (with X' reduced) that is a bijection on points is an isomorphism); moreover, a seminormal scheme which is  $G_1$  and  $S_2$  has only double crossing singularities in codimension 1, ie it is demi-normal. We refer to [Tra70, GT80, KSS10]) for the previous assertions. This leads to the following definition:

**Definition 2.8.** — We will say that X has semi-log canonical singularities if:

- 1. X is  $G_1$  and  $S_2$ ,
- 2.  $K_X$  is  $\mathbb{Q}$ -Cartier,
- 3. The pair  $(X^{\nu}, C_{X^{\nu}})$  is log-canonical.

If X has semi-log canonical singularities (slc), then  $C_{X^{\nu}}$  is necessarily reduced, and therefore the codimension 1 singularities of X are only double crossing as we explained above. This assumption is usually added in the definitions (cf [Kol, Kov12]), but we may keep it or not without any change. This justifies the seemingly different definition given in the previous section. Finally, we can give the definition of a stable variety:

# **Definition 2.9.** — We say that X is stable if

- 1. X is projective,
- 2. X has semi-log canonical singularities,
- 3.  $K_X$  is  $\mathbb{Q}$ -ample.

Singularities and Kähler-Einstein metrics: a summary. — If we take a closer look at the proof of Proposition 2.6, we see that we did not use all of the properties of a stable variety to see that a Kähler-Einstein metric always extend. Actually, we just used the fact that the conductor was a divisor. Therefore, using the existence and regularity results that we are going to prove in the next sections, and the restriction on the singularities of a pair carrying a Kähler-Einstein metric (cf Proposition 5.1), we can summarize the problem of the existence of a Kähler-Einstein metric on a stable variety in the following statement:

**Theorem 2.10.** — Let X be reduced n-equidimensional projective scheme of finite type over  $\mathbb{C}$ , satisfying the conditions  $G_1$  and  $S_2$ , and such that  $K_X$  is an ample  $\mathbb{Q}$ -line bundle. Then the following are equivalent:

- 1. There exists a Kähler form  $\omega$  on  $X_{\text{reg}}$  such that  $\operatorname{Ric} \omega = -\omega$  and  $\int_{X_{\text{reg}}} \omega^n = c_1(K_X)^n$ ,
- 2. There exists  $\omega$  as above which extends to define a positive current in  $c_1(K_X)$ ,
- 3. X has semi-log canonical singularities, ie X is stable.

Moreover, by the results of Odaka [Oda08, Oda11], the latter condition is equivalent to: 4. The pair  $(X, K_X)$  is K-stable.

### 3. Variational solutions

**3.1. General setting.** — Consider the following general setting: X is a compact Kähler manifold and  $[\omega]$  a big class, with  $\omega$  smooth (but not necessarily positive). We say that a function  $u \in PSH(X,\omega)$  has full Monge-Ampère mass, and we will write  $u \in \mathcal{E}(X,\omega)$ , if the total mass of MA(u) is equal to the volume of the class  $[\omega]$ , where the volume in question may be defined by  $V := \text{vol}([\omega]) := \int_X MA(u_{\min})$ , for  $u_{\min}$  any element in  $PSH(X,\omega)$  with minimal singularities, cf §1.

We now recall an important subspace of  $\mathcal{E}(X,\omega)$  denoted by  $\mathcal{E}^1(X,\omega)$ , and consisting of functions with finite energy. The energy  $\mathcal{E}(u)$  of an  $\omega$ -psh function u (not necessarily in  $\mathcal{E}(X,\omega)$ ) is defined in the following way (cf [GZ07, BEGZ10, BBGZ09] for more details – the energy is sometimes denoted by E in the aforementioned papers).

First, if  $u \in PSH(X, \omega)$  has minimal singularities, we set

$$\mathcal{E}(u) := \frac{1}{n+1} \sum_{j=0}^{n} \int_{X} (u - V_{\theta}) \, \mathrm{MA}(u^{(j)}, V_{\theta}^{(n-j)})$$

where MA is the mixed non-pluripolar Monge-Ampère operator. If now u is any  $\omega$ -psh function, we defined

$$\mathcal{E}(u) := \inf \{ \mathcal{E}(v) \mid v \in PSH(X, \omega) \text{ with minimal singularities, } v \geqslant u \}$$

Then we set  $\mathcal{E}^1(X,\omega) := \{u \in PSH(X,\omega), \mathcal{E}(u) > -\infty\}$ . Actually, [**BEGZ10**, Proposition 2.11] gives another characterization of this last space: a function  $u \in PSH(X,\omega)$  belongs to  $\mathcal{E}^1(X,\omega)$  if and only if  $u \in \mathcal{E}(X,\omega)$  and  $\int_X (u - V_\theta) MA(u) < +\infty$  (and for any  $u \in \mathcal{E}(X,\omega)$ , the explicit integral formula for  $\mathcal{E}(u)$  above is still valid). Using this result, it becomes clear that  $\mathcal{E}^1(X,\omega) \subset \mathcal{E}(X,\omega)$  as announced.

We should finally add that  $\mathcal{E}$  is an upper-semicontinuous (usc) concave functional on  $PSH(X,\omega)$ , and that it is the normalized primitive of the Monge-Ampère operator, i.e.

(3.1) 
$$(d\mathcal{E})_u = \frac{1}{V} \operatorname{MA}(u)^n.$$

**3.2.** Uniqueness. — Given a measure  $\mu$  on X (possible non-finite) we consider the following MA-equation for  $u \in PSH(X, \omega)$  attached to the pair  $(\omega, \mu)$ :

$$\omega_u^n = e^u \mu,$$

where  $\omega_u^n := \mathrm{MA}(u)$  is the non-pluripolar Monge-Ampère operator as before. This equation is equivalent to the following normalized MA-equation on  $\mathcal{E}(X,\omega)/\mathbb{R}$ :

(3.3) 
$$\frac{\omega_u^n}{V} = \frac{e^u \mu}{\int e^u \mu},$$

The equivalence follows immediately from the  $\mathbb{R}$ -invariance of the latter equation and the substitution  $u \mapsto u - \log \int e^u \mu$  which maps solutions of equation (3.2) to solutions of the equation (3.3).

**Proposition 3.1.** — Any two solutions u and v of the MA-equation (3.2) such that u and v are in  $\mathcal{E}(X)$  coincide.

*Proof.* — This is an immediate consequence of the comparison principle [**BEGZ10**, Corollary 2.3]: if u and v are in  $\mathcal{E}(X)$  then

$$\int_{\{u < v\}} \mathrm{MA}(v) \leqslant \int_{\{u < v\}} \mathrm{MA}(u)$$

But the MA above then forces u = v a.e wrt the measure  $\mu$ . Since  $\mu$  cannot charge pluripolar sets (as MA(u) does not) it follows that u = v away from a pluripolar set and hence everywhere, by basic properties of psh functions.

**3.3. Existence results for log canonical pairs.** — Let (X, D) be a log canonical pair such that the log canonical divisor  $K_X + D$  is big. Assume that (X, D) is a log smooth, i.e. X is smooth and

$$D = \sum_{i} c_i D_i$$

is a normal crossings divisor with  $c_i \in ]-\infty,1]$ . To the pair (X,D) we can associate the following Kähler-Einstein type equation for a metric  $\phi$  on  $L:=K_X+D$ :

$$(3.4) (dd^c \phi)^n = e^{\phi - \phi_D},$$

where  $\phi_D = \sum_i c_i \log |s_i|^2$  and  $s_i$  are sections cutting out the divisors  $D_i$  above.

**Theorem 3.2.** — There is a unique finite energy solution  $\phi$  to the equation above.

*Proof.* — The proof is a modification of the variational approach in [BBGZ09] (concerning the case when D is trivial). To explain this we fix a smooth form  $\omega \in c_1(K_X + D)$ . Then the equation above is equivalent to a Monge-Ampère equation for an  $\omega$ -psh function u:

$$\omega_u^n = e^u \mu$$

where the measure  $\mu$  is of the form  $\mu = \rho dV$  for a function  $\rho$  in  $L^{1-\delta}(X)$  (but  $\rho$  is not in  $L^1(X)$ !). We let

$$\mathcal{L}(u) := -\log \int e^u \mu$$

Then, at least formally, solutions of equation (3.5) are critical points of the functional

$$\mathcal{G}(u) := \mathcal{E}(u) + \mathcal{L}(u).$$

in view of the equation (3.1) satisfied by  $\mathcal{E}$ .  $\mathcal{L}$  also defines an usc concave functional on  $PSH(X,\omega)$  and we let  $\mathcal{L}(X,\omega) := \{\mathcal{L} > -\infty\}$  (the upper semi-continuity follows from Fatou's lemma).

Note that Lemma 3.3 below guarantees that the intersection  $\mathcal{E}^1(X,\omega) \cap \mathcal{L}(X,\omega)$  is nonempty. Hence,  $\mathcal{G}(u)$  is not identically equal to  $-\infty$  on its domain of definition that we will take to be  $\mathcal{E}^1(X,\omega)$  (equipped with the usual  $L^1(X)$ -topology).

Next, we observe that

(3.6) 
$$\mathcal{G}(u) \leqslant \mathcal{E}(u) - \int u\mu_0 + C''$$

Indeed, since  $\mu \geqslant C\mu_0$ , where  $\mu_0$  is finite measure on X integrating all quasi-psh functions on X (in our case we may take  $\mu_0 = ||s'|| dV$  for some holomorphic section s' defined by the negative coefficients of D):

$$\int e^u \mu \geqslant C \int e^u \mu_0$$

and hence

$$\mathcal{L}(u) \leqslant C' - \log \int e^u \mu_0 \leqslant C'' - \int u \mu_0$$

using Jensen's inequality, which proves (3.6). In particular,  $\mathcal{G}(u)$  is bounded from above. Indeed, by scaling invariance we may assume that  $\sup_X u = 0$  and then use that, by basic compactness properties of  $\omega$ -psh functions,  $\sup u \leq \int u\mu_0 + C$ .

Let now  $u_i \in \mathcal{E}^1(X,\omega)$  be a sequence such that

$$\mathcal{G}(u_j) \to \sup_{\mathcal{E}^1(X,\omega)} \mathcal{G} := S < \infty$$

Again, by the scale invariance of  $\mathcal{G}$  we may assume that  $\sup_X u_i = 0$ . In particular,

$$\mathcal{L}(u_j) \geqslant S/2 - \mathcal{E}(u_j)$$

for  $j > j_0$ . But, by (3.6),  $\mathcal{E}(u_j)$  is bounded from below and hence there is a constant C such that

$$\mathcal{E}(u_i) \geqslant -C, \ \mathcal{L}(u_i) \geqslant -C$$

Let now  $u_*$  be a limit point of  $u_j$ . By upper-semicontinuity we have that

$$\mathcal{E}(u) \geqslant -C, \ \mathcal{L}(u) \geqslant -C$$

Finally, we note that  $u_*$  satisfies the equation (3.5) by applying the projection argument from [BBGZ09] as follows. Fixing  $v \in \mathcal{C}^{\infty}(X)$  let  $f(t) := \mathcal{E}_{\omega}(P_{\omega}(u_* + tv) + \mathcal{L}(u_* + tv))$ , where

$$P_{\omega}(u)(x) := \sup\{v(x) : v \leqslant u, v \in PSH(X, \omega)\}$$

(note that f(t) is finite for any t). The functional  $\mathcal{L}(u)$  is decreasing in u and hence the sup of f(t) on  $\mathbb{R}$  is attained for t = 0. Now  $\mathcal{E}_{\omega} \circ P_{\omega}$  is differentiable with differential  $MA(P_{\omega}u)$  at u [BBGZ09]. Hence, the condition df/dt = 0 for t = 0 gives that the variational equation (3.5) holds when integrated against any  $v \in \mathcal{C}^{\infty}(X)$ .

Let us now prove the following result, that we used in the proof:

**Lemma 3.3.** — Let (X,D) be a log smooth pair such that  $L := K_X + D$  is big. Let  $\theta$  be a smooth (1,1) form whose cohomology class is  $c_1(L)$ . Let  $s_0$  be a section of D, and  $|\cdot|$  a smooth hermitian metric on  $\mathcal{O}_X(D)$ . Then there exists a  $\theta$ -psh function  $u \in \mathcal{E}^1(X,\theta)$  such that  $e^u/|s_0|^2$  is integrable.

*Proof.* — As L is big, the exists m big enough such that mL - D is effective. We choose t a holomorphic section of mL - D, and consider  $s := s_0 t$  which is a section of mL vanishing along D. Let  $h_0$  be an smooth hermitian metric on L with curvature form  $\theta$ , and let  $V_{\theta}$  be the upper envelope of all (normalized)  $\theta$ -psh functions. We define on mL the hermitian metric  $h := h_0^{\otimes m} e^{-mV_{\theta}}$ . For  $0 < \alpha < 1$  small enough we claim that the function

$$u := V_{\theta} - \left(-\frac{1}{m}\log|s|_{h}^{2}\right)^{\alpha}$$

suits our requirements.

First of all, it is  $\theta$ -psh because of the general following fact: if  $\psi$  is  $\theta$ -psh and  $\chi : \mathbb{R} \to \mathbb{R}$  is convex and non-decreasing satisfying  $\chi' \leq 1$ , then  $V_{\theta} + \chi(\psi - V_{\theta})$  is  $\theta$ -psh. Indeed,  $dd^{c}(V_{\theta} + \chi(\psi - V_{\theta})) = (1 - \chi')(\theta + dd^{c}V_{\theta}) + \chi'(\theta + dd^{c}\psi) + \chi''|d(\psi - V_{\theta})|^{2} \text{ where } \chi' \text{ and } \chi''$ 

are evaluated at  $\psi - V_{\theta}$ . Now we apply this to  $\psi = 1/m \log |s|_{h_0^m}^2$ .

For the integrability property, we use the following inequality for x a real number (big enough):  $x^{\alpha} \ge (n+1) \log x - C$  for some C > 0 depending only on  $\alpha$ . Now we observe that  $V_{\theta} + \chi(\psi - V_{\theta}) \le \chi(\psi)$ : indeed,  $\chi(\psi - V_{\theta}) - \chi(\psi) \le \sup \chi' \cdot (-V_{\theta}) \le -V_{\theta}$ , so that in our case,  $u \le (-\frac{1}{m} \log |s|_{h_0}^2)^{\alpha}$ . If we apply the basic inequality stated above to  $x = -\frac{1}{m} \log |s|_{h_0}^2$  which can be made big enough by multiplying  $h_0$  by a big constant (this does not change the curvature), we get

$$e^u \leqslant C \left( -\frac{1}{m} \log |s|_{h_0}^2 \right)^{-(n+1)}$$

As D has snc support, and |t| is bounded from above, we are left to check that the integral

$$\int_D \frac{dV}{\prod_{i \leqslant n} |z_i|^2 \cdot \log^{n+1}(\prod_{i \leqslant n} |z_i|^2)}$$

over the unit polydisc D in  $\mathbb{C}^n$  converges. But after a polar change of coordinate, we are led to estimate  $\int_{[0,1]^n} \frac{dx_1 \cdots dx_n}{\prod_{i \leqslant n} x_i \cdot \log^{n+1}(\prod_{i \leqslant n} x_i^2)}$ , which equals  $\frac{1}{2^{n+1}n} \int_{[0,1]^{n-1}} \frac{dx_1 \cdots dx_{n-1}}{\prod_{i \leqslant n-1} x_i \cdot \log^n(\prod_{i \leqslant n-1} x_i)}$ . By induction, and using the Poincaré case, it concludes.

Finally, one has to check that  $u \in \mathcal{E}^1(X,\theta)$ . We compute the capacity  $\operatorname{Cap}_{\theta}(u < V_{\theta} - t)$  for t big. But  $(u < V_{\theta} - t) = \left(\frac{1}{m} \log |s|_h^2 < -t^{1/\alpha}\right) \subset \left(\left(\frac{1}{m} \log |s|_{h_0}^2 < -t^{1/\alpha}\right)\right)$  and thus  $\operatorname{Cap}_{\theta}(u < V_{\theta} - t) \leqslant \frac{C}{t^{1/\alpha}}$  because for every  $\theta$ -psh function  $\psi$ , one has  $\operatorname{Cap}_{\theta}(\psi < -t) \leqslant \frac{C_{\psi}}{t}$  (this is an easy generalization of [GZ05, Proposition 2.6]). Therefore, if  $\alpha < \frac{1}{n+1}$ , one has

$$\int_0^{+\infty} t^n \operatorname{Cap}_{\theta}(u < V_{\theta} - t) dt < +\infty$$

which, using the characterization given in [BBGZ09, Lemma 2.9], ends the proof of the lemma.

**3.4. Stability under perturbations.** — Let now  $L_j$  be sequence of semi-positive line bundles such that there exists forms  $\omega_j \in c_1(L_j)$  and  $\omega \in c_1(L)$  satisfying  $\omega_j \geq 0$  and  $\omega_j \to \omega$  in the  $C^{\infty}$ -topology. Let  $\mu_j := \rho_j dV$  be a sequence of measures on X such that  $\rho_j$  increases to  $\rho$  (= the density defined in the previous section) and assume moreover that  $\mu_j \geq C\mu_0$  for some finite measure  $\mu_0$  as in the previous section. Consider the following Monge-Ampère equations for  $u_j \in \mathcal{E}(X, \omega_j)$  (and sup-normalized):

$$\omega_{u_j}^n/V = \frac{e^{u_j}\mu_j}{\int_X e^{u_j}\mu_j}$$

and similarly

$$\omega_u^n/V = \frac{e^u \mu}{\int_X e^u \mu}$$

for  $u \in \mathcal{E}(X, \omega)$ .

**Theorem 3.4.** — The unique sup-normalized solution  $u_j$  of the first equation above converges, in the  $L^1(X)$ -topology, to the unique sup-normalized solution u to the the latter equation. Equivalently, the solutions  $v_j$  of the corresponding non-normalized equations converge in  $L^1(X)$  to v solving the corresponding limiting non-normalized equation.

*Proof.* — First observe that, by assumption,  $u \in \mathcal{E}(X, \omega_j)$  for any j and moreover

$$(3.7) -C \leqslant \mathcal{G}_i(u^j) < \infty$$

where  $\mathcal{G}_j$  denotes the functional determined by the pair  $(\omega_j, \mu_j)$ . Indeed, by (ii) above we have

$$-\mathcal{E}_{\omega_j}(u) := \int u_s \sum_{k=0}^n \omega_u^k \wedge \omega_j^{n-k} \geqslant -C'$$

using that  $\mathcal{E}_{\omega}(u) > -\infty$  and  $\omega_j \to \omega$  in the  $\mathcal{C}^{\infty}$ -topology (compare the proof of the lemma below) and the claim (3.7) then follows from

$$\log \int e^u \mu_j \leqslant C + \log \int e^u \mu < C'$$

(since, by assumption  $\mu_j \leqslant \mu$ ). Let now  $u_j$  be the sup-normalized maximizer of  $\mathcal{G}_j$  so that

$$\mathcal{G}_j(u_j) \geqslant \mathcal{G}_j(u) \geqslant -C$$

Using the assumption  $\mu_i \geqslant C\mu_0$  and Jensen's inequality as before we also have

$$\mathcal{G}_j(u) \leqslant \mathcal{E}_{\omega_j}(u) - \int_X u\mu_0 + C''$$

Next, the sequence  $u_j$  enjoys the usual compactness properties of quasi-psh functions and in particular,

$$0 = \sup u_j \leqslant \int u_j \mu_0 + C_\theta$$

Combining the latter two inequalities we deduce

$$\mathcal{E}_{\omega_i}(u_i) \geqslant -C$$

Let now u be an  $L^1$ -limit point in  $PSH(X, \omega)$  of the sequence  $u_j$ . By the previous inequality and the lemma below we have

$$\mathcal{E}_{\omega}(u) \geqslant \limsup_{j} \mathcal{E}_{\omega_{j}}(u_{j}) \geqslant -C$$

Moreover, by Fatou's lemma

$$\mathcal{L}(u) \geqslant \limsup_{j} \mathcal{L}_{j}(u_{j})$$

Hence,

$$\mathcal{G}(u) \geqslant \limsup_{i} \mathcal{G}_{j}(u_{j})$$

Next, since by assumption  $\omega_j \geqslant \omega$  we have that  $PSH(X,\omega) \subset PSH(X,\omega_j)$  and hence any given  $v \in PSH(X,\omega)$  is a candidate for the sup of  $\mathcal{G}_j$  so that

$$\mathcal{G}(u) \geqslant \limsup_{j} \mathcal{G}_{j}(u_{j}) \geqslant \limsup_{j} \mathcal{G}_{j}(v) = \mathcal{G}(v)$$

and since v was an arbitrary element in  $PSH(X, \omega)$  we then conclude that u is a maximizer of  $\mathcal{G}$ , as desired.

Finally, we have  $v_j = u_j + \mathcal{L}_j(u_j)$  and as shown above  $\mathcal{G}_j(u_j)$  is a bounded sequence (more precisely, it converges to the maximal value S of  $\mathcal{G}$ ) and  $0 \le -\mathcal{E}(u_j) \le C$ , which implies that  $\mathcal{L}_j(u_j)$  is also a bounded sequence. After passing to a subsequence we may thus assume that  $\mathcal{L}_j(u_j) \to l \in \mathbb{R}$  so that  $v_j \to v := u + l$ , solving the desired equation (and  $v \in \mathcal{E}^1(X, \omega)$ ). By the uniqueness of solutions of the latter equation this means that the whole sequence  $u_j$  converges to v, which concludes the proof.

**Lemma 3.5.** — Let  $[\omega_j]$  and  $[\omega]$  be semi-positive big classes such that  $\omega_j \to \omega$  in the  $C^{\infty}$ -topology (without making any semi-positive assumptions on the smooth forms  $\omega_j$  and  $\omega$ ). If  $u_j \in \mathcal{E}(X, \omega_j)$  (and  $u \in \mathcal{E}(X, \omega)$ ) such that  $u_j \to u$  in  $L^1(X)$ , then

$$\mathcal{E}_{\omega}(u) \geqslant \limsup_{j} \mathcal{E}_{\omega_{j}}(u_{j})$$

*Proof.* — When  $\omega_j = \omega$  the lemma amounts to the well-known fact that  $\mathcal{E}_{\omega}$  is usc. We may as well assume that  $\sup u_j = \sup u = 0$ .

First of all if  $\limsup_j \mathcal{E}_{\omega_j}(u_j) = -\infty$  then the statement is trivial and hence we may as well in ther proof assume (after perhaps reindexing the sequence) that  $\mathcal{E}_{\omega_j}(u_j) \geqslant -C$ . Fix  $\varepsilon$  and A above so that  $PSH(X, \omega_j)$  and  $PSH(X, \omega_j)$  are both included in  $PSH(X, \omega_j + \varepsilon \omega_A)$ . Since  $\mathcal{E}_{\omega + \varepsilon \omega_A}$  is usc we have

$$\mathcal{E}_{\omega+\varepsilon\omega_A}(u) \geqslant \limsup_{i} \mathcal{E}_{\omega+\varepsilon\omega_A}(u_i)$$

and letting  $\varepsilon \to 0$  thus gives

$$\mathcal{E}_{\omega}(u) \geqslant \limsup_{j} \mathcal{E}_{\omega + \varepsilon \omega_{A}}(u_{j})$$

Next, for j sufficiently large we have  $\omega \leq \omega_j + \varepsilon \omega_A$ . Hence,

$$-\mathcal{E}_{\omega+\varepsilon\omega_A}(u_j) = \sum_k \int (-u_j)(\omega + dd^c u_j + \varepsilon\omega_A)^k \wedge (\omega + \varepsilon\omega_A)^{n-k} \leqslant$$

$$\sum_k \int (-u_j)(\omega_j + dd^c u_j + 2\varepsilon\omega_A)^k \wedge (\omega_j + 2\varepsilon\omega_A)^{n-k} =$$

$$= -\mathcal{E}_{\omega_j}(u_j) + \varepsilon \int_X \sum_k c_k (-u_j)(\omega_A)^k \wedge (\omega_{u_j})^{n-k}$$

Finally, it is well-known [BBGZ09] that if  $-\mathcal{E}_{\omega_j}(u_j) \leq C$  and  $\sup u_j = 0$ , then the integral in the rhs above is bounded by a constant C' and hence

$$\mathcal{E}_{\omega}(u) \geqslant \limsup_{j} \mathcal{E}_{\omega_{j}}(u_{j}) - C' \varepsilon$$

so that letting  $\varepsilon \to 0$  concludes the proof.

Corollary 3.6. — Let (X, D) be a log smooth log canonical pair (in particular, the coefficients of D are in  $]-\infty,1]$ ) and assume that  $L:=K_X+D$  is semi-positive and big. Fixing an ample line bundle A set  $L_j:=L+A/j$ . Let  $\psi_j$  be a decreasing sequence of smooth metrics on on the klt part  $D_{klt}$  of D such that  $\psi_j \to \phi_{klt}$  (where  $dd^c\phi_{klt}=[D_{klt}]$ ) and consider the Monge-Ampère equations

$$(dd^c\phi_i)^n = e^{\phi_j - \psi_j - \phi_D}$$

for a metric  $\phi_j \in \mathcal{E}(X, L_j)$ . The equations admit unique solutions  $\phi_j$  and moreover  $\phi_j \to \phi_{KE}$  in  $L^1$  where  $\phi_{KE}$  is the unique solution of equation (3.4).

## 4. Smoothness of the Kähler-Einstein metric

Before we go into the details of the proof of the regularity theorem, we would like to give an overview of previous related results and underline the main differences that will appear in our specific case, namely the case of general log canonical pairs. As we will rely on the so called logarithmic case (ie (X, D) is log smooth, D is reduced, and  $K_X + D$  is ample), the next section will be devoted to recall some of the main tools appearing in this setting. Then, we will give the proof of the main regularity theorem, which will constitute the core of this section.

**4.1. Special features in the log canonical case.** — We should first mention the case of varieties with log terminal singularities, or more generally klt pairs, which correspond to the pairs where the discrepancies  $a_i$  defined earlier satisfy the strict inequalities  $a_i > -1$ . Then the situation is relatively well understood: In the non-positively curved case, we know that the Kähler-Einstein metric exists, is unique, has bounded potential, and induces on the regular locus a genuine Kähler-Einstein metric (see e.g. [BEGZ10, EGZ08, EGZ09, DP10]). As for the case of positive curvature, or log Fano manifolds, then there exist criteria (like the properness of the Mabuchi functional) to guarantee the existence and uniqueness (modulo automorphisms of X) of a Kähler-Einstein metric [BBE+11]; this metric is also known to have bounded potential and to be smooth on the regular locus of X (see again [BBE+11, Pău08]).

However, the behavior of the Kähler-Einstein metric near the singularities of X is mostly unknown (except if the singularities are orbifold). In the case of a klt pair, we know that the metric will not be smooth along the divisor, but its singularities can sometimes be understood outside of the singular part of (X, D). For example, a recent result in this direction states (under some technical assumption) that the Kähler-Einstein metric has cone singularities near each point where (X, D) is log smooth, ie X is smooth and D has simple normal crossing support (cf [Gue12a]).

When (X, D) is now a log smooth pair, the situation gets easier because there is no more loss of positivity coming from the resolution of singularities. For example, if the coefficients

of D are in [0,1) (the pair is then klt), the Kähler-Einstein metric is known to have cone singularities along D at least under some assumptions like D is smooth [JMR11], the coefficients are greater than 1/2 [CGP11], or both [Bre11]. The general case has been announced in [MR12].

When now every coefficient of D is equal to 1, and  $K_X + D$  is ample, then we know from the work of Kobayashi [Kob84] and Tian-Yau [TY87] that there exists a unique complete Kähler-Einstein metric having Poincaré singularities along D. The situation where the coefficients of D are in [1/2, 1] behaves like a product of cone and Poincaré geometries and was studied in [Gue12b].

In a slightly different direction, Tsuji [Tsu88a] has considered the case of a singular variety with ample canonical line bundle such that only one divisor appears in its resolution, with discrepancy equal to -1. Finally, Wu [Wu08, Wu09] has worked out the case of a quasi-projective manifold compactified by a snc divisor  $\sum D_i$  such that  $K_X + \sum a_i D_i$  is ample for some coefficients  $a_i \ge -1$ . In our case however, such a strong positivity assumption will never happen as soon as we have to perform a non-trivial resolution.

As one can already observe in the log smooth case studied by Kobayashi and Tian-Yau, the log canonical case is very different from the klt case. Let us mention some striking divergences: first of all, the potentials are no more bounded even in the ample case so that the solution does not have minimal singularities. Moreover, the Kähler-Einstein equation in this setting is closely related to a negative curvature geometry. Indeed, if we first consider the Ricci-flat case, then it is impossible to write the equation on the whole X. Indeed, the current solution obtained on  $X_{\text{reg}}$  will not have finite mass near the singularities, and hence it will not extend as a global positive current on X. This phenomenon already happens in  $[\mathbf{TY90}]$ . Finally, it has been proven in  $[\mathbf{BBE+11}]$ , Proposition 3.8] that any pair (X, D) with X normal and  $-(K_X + D)$  ample admitting a Kähler metric  $\omega$  on  $X_{\text{reg}}$  with continuous potentials solution of  $\text{Ric }\omega = \omega + [D]$  is necessarily klt. Therefore it is pointless to look for positively curved Kähler-Einstein metric in the general setting of log canonical pairs instead of klt pairs.

To finish this discussion, let us stress the fact that the class of varieties with semi-log canonical singularities can be realized as a subclass of log canonical pairs (cf definition 2.8). This is the largest "reasonable" class to look for Kähler-Einstein metrics: for example, if X is a smooth Fano manifold carrying a smooth divisor  $D \in |-K_X|$ , then for any  $\varepsilon > 0$ , one has  $K_X + (1 + \varepsilon)D > 0$ ; however, there is no smooth Kähler-Einstein metric with negative scalar curvature on  $X \setminus D$ . Indeed, its existence would contradict the Yau-Schwarz lemma applied with the complete Ricci-flat Kähler metric constructed in [**TY90**].

Moreover, we will see that the existence of a negatively curved Kähler-Einstein metric on the regular part of a normal projective variety with maximal volume forces the singularities to be at worst log canonical, cf. Proposition 5.1. **4.2.** The logarithmic case. — In this section, we will briefly recall the Theorem of Kobayashi and Tian-Yau constructing negatively-curved Kähler-Einstein metrics on quasi-projective varieties  $X \setminus D$  where D is a reduced divisor with simple normal crossings, and  $K_X + D$  is ample. In the course of the proof of Theorem 4.6, we will use in an essential manner the functional spaces introduced by these authors, namely the "quasi-coordinates" version of the usual Hölder spaces  $\mathscr{C}^{k,\alpha}$ . For now,  $X_0$  will denote  $X \setminus D$ .

**Definition 4.1.** — We say that a Kähler metric  $\omega$  on  $X_0$  is of Carlson-Griffiths type if there exists a Kähler form  $\omega_0$  on X such that  $\omega = \omega_0 - \sum_K dd^c \log \log \frac{1}{|s_k|^2}$ .

In [CG72], Carlson and Griffiths introduced such a metric for some  $\omega_0 \in c_1(K_X + D)$ , but one can easily see that such a metric always exists on a Kähler manifold without assumptions on the bundle  $K_X + D$ . One can also observe that the existence of such a metric  $\omega$  forces the cohomology class  $\{\omega\}$  to be Kähler by Demailly's regularization theorem [Dem82, Dem92].

The reason why we exhibit this particular class of Kähler metrics on  $X_0$  having Poincaré singularities along D is that we have an exact knowledge on its behaviour along D, which is much more precise that knowing its membership in the aforementioned class. This is precisely the class in which one will look for a Kähler-Einstein metric, so that one needs to define the appropriate analogue of the usual Hölder spaces  $\mathscr{C}^{k,\alpha}$ . And to do so, one may (almost) boil down to the usual euclidian situation.

The key point is that  $(X_0, \omega)$  has bounded geometry at any order. Let us get a bit more into the details. To simplify the notations, we will assume that D is irreducible, so that locally near a point of D,  $X_0$  is biholomorphic to  $\mathbb{D}^* \times \mathbb{D}^{n-1}$ , where  $\mathbb{D}$  (resp.  $\mathbb{D}^*$ ) is the unit disc (resp. punctured disc) of  $\mathbb{C}$ . We want to show that, roughly speaking, the components of  $\omega$  in some appropriate coordinates have bounded derivatives at any order. The right way to formalize it consists in introducing quasi-coordinates: they are maps from an open subset  $V \subset \mathbb{C}^n$  to  $X_0$  having maximal rank everywhere. So they are just locally invertible, but these maps are not injective in general.

To construct such quasi-coordinates on  $X_0$ , we start from the univeral covering map  $\pi: \mathbb{D} \to \mathbb{D}^*$ , given by  $\pi(w) = e^{\frac{w+1}{w-1}}$ . Formally, it sends 1 to 0. The idea is to restrict  $\pi$  to some fixed ball B(0,R) with 1/2 < R < 1, and compose it (at the source) with a biholomorphism  $\Phi_{\eta}$  of  $\mathbb{D}$  sending 0 to  $\eta$ , where  $\eta$  is a real parameter which we will take close to 1. If one wants to write a formula, we set  $\Phi_{\eta}(w) = \frac{w+\eta}{1+\eta w}$ , so that the quasi-coordinate maps are given by  $\Psi_{\eta} = \pi \circ \Phi_{\eta} \times \mathrm{Id}_{\mathbb{D}^{n-1}}: V = B(0,R) \times \mathbb{D}^{n-1} \to \mathbb{D}^*$ , with  $\Psi_{\eta}(v,v_2,\ldots,v_n) = (e^{\frac{1+\eta}{1-\eta}\frac{v+1}{v-1}},v_2,\ldots,v_n)$ .

Once we have said this, it is easy to see that  $X_0$  is covered by the images  $\Psi_{\eta}(V)$  when  $\eta$  goes to 1, and for all the trivializing charts for X, which are in finite number. Now, an easy computation shows that the derivatives of the components of  $\omega$  with respect to the  $v_i$ 's are bounded uniformly in  $\eta$ . This can be thought as a consequence of the fact that the Poincaré metric is invariant by any biholomorphism of the disc.

At this point, it is natural to introduce the Hölder space of  $\mathscr{C}_{qc}^{k,\alpha}$ -functions on  $X_0$  using the previously introduced quasi-coordinates:

**Definition 4.2.** — For a non-negative integer k, a real number  $\alpha \in ]0.1[$ , we define:

$$\mathscr{C}_{qc}^{k,\alpha}(X_0) = \{ u \in \mathscr{C}^k(X_0); \sup_{V,n} ||u \circ \Psi_\eta||_{k,\alpha} < +\infty \}$$

where the supremum is taken over all our quasi-coordinate maps V (which cover  $X_0$ ). Here  $||\cdot||_{k,\alpha}$  denotes the standard  $\mathscr{C}_{qc}^{k,\alpha}$ -norm for functions defined on a open subset of  $\mathbb{C}^n$ .

The following fact, though easy, is very important (see e.g [Kob84] or [Gue12b, Lemma 1.6] for a detailed proof):

**Lemma 4.3.** — Let  $\omega$  be a Carlson-Griffiths type metric on  $X_0$ , and  $\omega_0$  some Kähler metric on X. Then

$$F_0 := \log \left( \prod |s_k|^2 \log^2 |s_k|^2 \cdot \omega^n / \omega_0^n \right)$$

belongs to the space  $\mathscr{C}^{k,\alpha}_{qc}(X_0)$  for every k and  $\alpha$ .

Finding the Kähler-Einstein metric consists then in showing that the Monge-Ampère equation  $(\omega + dd^c\varphi)^n = e^{\varphi + f}\omega^n$  has a unique solution  $\varphi \in \mathscr{C}^{k,\alpha}_{qc}(X_0)$  for all functions  $f \in \mathscr{C}^{k,\alpha}_{qc}(X_0)$  with  $k \geqslant 3$ . This can be done using the continuity method in the quasi-coordinates. In particular, applying this to  $f = F := -\log\left(\prod |s_k|^2 \log^2 |s_k|^2 \cdot \omega^n/\omega_0^n\right) + (\text{smooth terms on } X)$ , which the previous lemma allows to do, this will prove the existence of a negatively curved Kähler-Einstein metric, which is equivalent to  $\omega$  (in the *strong* sense:  $\varphi \in \mathscr{C}^{k,\alpha}_{qc}(X_0)$  for all  $k,\alpha$ ).

In this continuity method, one needs to obtain first uniform estimates; they follow from a consequence of Yau's maximum principle for complete manifolds which we recall here (see [CY80, Proposition 4.1]):

**Proposition 4.4.** — Let  $(X,\omega)$  be a n-dimensional complete Kähler manifold, and  $F \in \mathscr{C}^2(X)$  bounded from above. We assume that we are given  $u \in \mathscr{C}^2(X)$  satisfying  $\omega + dd^c u > 0$  and

$$(\omega + dd^c u)^n = e^{u+F} \omega^n$$

Suppose that the bisectional curvature of  $(X, \omega)$  is bounded below by some constant, and that u is bounded from below. Then

$$\inf_X u \geqslant -\sup_X F \quad and \quad \sup_X u \leqslant -\inf_X F$$

There are similar results for the Laplacian estimates, but as we will not use them directly, we do not state them here. To summarize the discussion, one obtains:

Theorem 4.5 (Kobayashi [Kob84], Tian-Yau [TY87]). — Let X be a compact Kähler manifold, D a reduced divisor with simple normal crossings,  $\omega$  a Kähler form of Carlson-Griffiths type on  $X \setminus D$ , and  $F \in \mathscr{C}^{k,\alpha}_{qc}(X \setminus D)$  for some  $k \geqslant 3$ . Then there exists  $\varphi \in \mathscr{C}^{k,\alpha}_{qc}(X \setminus D)$  solution to the following equation:

$$(\omega + dd^c \varphi)^n = e^{\varphi + F} \omega^n$$

In particular if  $K_X + D$  is ample, then there exists a (unique) Kähler-Einstein metric of curvature -1 equivalent to  $\omega$ .

**4.3. Statement of the regularity theorem.** — In this section, we prove that the Kähler-Einstein metric attached to a log canonical pair (X, D) (satisfying  $K_X + D$  ample) by Theorem 3.2 is smooth on  $X_0 = X_{\text{reg}} \setminus \text{Supp}(D)$ . As usual, we will work on a log resolution  $\pi: X' \to X$ , where:

$$K_{X'} = \pi^*(K_X + D) + \sum a_i E_i$$

 $E_i$  being either an exceptional divisor or a component of the strict transform of D, and the coefficients  $a_i$  (called discrepencies) satisfy the inequality  $a_i \ge -1$ .

The Kähler-Einstein metric is given on X' by a (singular) psh weight  $\phi$  on  $\pi^*(K_X + D)$  satisfying

$$(dd^c\phi)^n = e^{\phi + \sum a_i \phi_{E_i}}$$

where  $\phi_{E_i}$  is a psh weight on  $\mathcal{O}_{X'}(E_i)$  such that  $dd^c\phi_{E_i}=[E_i]$ . So if in local coordinates,  $E_i$  is given by  $\{z_n=0\}$ , then  $\phi_{E_i}=\log|z_n|^2$ .

Our aim is to obtain regularity properties for the solutions of degenerate Monge-Ampère equations like the previous one; this is the content of the following theorem:

**Theorem 4.6.** — Let X be a compact Kähler manifold of dimension n, dV some volume form,  $D = \sum a_i D_i$  a  $\mathbb{R}$ -divisor with coefficients in  $(-\infty, 1]$  and defining sections  $s_i$ ,  $E = \sum c_{\alpha} E_{\alpha}$  an effective  $\mathbb{R}$ -divisor such that  $D_{red} + E$  has snc support, and  $\theta$  a semipositive form with  $\int_X \theta^n > 0$  such that  $\{\theta\} - c_1(E)$  is a Kähler class. Then the  $\theta$ -psh function  $\varphi$  with full Monge-Ampère mass, which is a solution of

$$\langle (\theta + dd^c \phi)^n \rangle = \frac{e^{\varphi} dV}{\prod_i |s_i|^{2a_j}}$$

is smooth outside of  $Supp(D) \cup Supp(E)$ .

Note that although  $\varphi$  has full Monge-Ampère mass, it is in general far from having minimal singularities as soon as some coefficient  $a_i$  of D equals 1. Think for example of the logarithmic case (a log smooth pair (X, D) where  $K_X + D$  is ample; then the potential of the Kähler-Einstein metric is not bounded whereas the class is ample.

Let us go back to the general Kähler-Einstein case. We would like to apply the previous results with E being some positive combination of the  $E_i$ 's. The problem is that there might

be no such divisors; for example if  $\pi$  happens to be a small resolution, its exceptional locus has codimension at least 2. Therefore we need to perform another modification.

On X',  $\pi^*(K_X + D)$  is no more ample, and by [BBP10, Proposition 1.5], its augmented base locus is  $\mathbb{B}_+(\pi^*(K_X + D)) = \pi^{-1}(\mathbb{B}_+(K_X + D)) \cup \operatorname{Exc}(\pi) = \operatorname{Exc}(\pi)$ , and lies above  $X_{\operatorname{sing}} \cup \operatorname{Supp}(D)$ . It is well-known that one can find a log resolution  $\mu: X'' \to X'$  of  $(X', \mathbb{B}_+(\pi^*(K_X + D)))$ , and an effective  $\mathbb{Q}$ -divisor F with snc support lying above  $\mathbb{B}_+(\pi^*(K_X + D))$  and such that  $\mu^*\pi^*(K_X + D) - F$  is ample. Moreover one can also assume that  $F + \sum E_i'$  has snc support, where  $E_i'$  denotes the strict transform of  $E_i$  by  $\mu$ .

Let us recall the argument. We start by resolving the singularities of a Kähler current  $T \geqslant \omega$  ( $\omega$  a Kähler form on X') in  $\pi^*(K_X + D)$  computing  $\mathbb{B}_+(\pi^*(K_X + D))$ , then we write Siu's decomposition  $\mu^*T = \theta + [D]$  with  $\theta$  semi-positive dominating  $\mu^*\omega$ , and D lying above  $\mathbb{B}_+(\pi^*(K_X + D))$ . Finally, we choose a  $\mu$ -exceptional  $\mathbb{Q}$ -divisor G such that -G is  $\mu$ -ample; it exists because  $\mu$  is a finite composition of blow-ups with smooth centers. Then it becomes clear that for  $\varepsilon > 0$  small enough,  $\{\mu^*\theta\} - \varepsilon G$  is a Kähler class, and we have  $\mu^*\pi^*(K_X + D) = (\{\mu^*\theta\} - \varepsilon G) + (\varepsilon G + D)$ , with  $\varepsilon G + D$  lying above  $\mathbb{B}_+(\pi^*(K_X + D))$  and having simple normal crossing support. If one had chosen a log resolution of the ideal sheaf generated by the augmented base locus of  $\pi^*(K_X + D)$  and the  $\mathcal{O}'_X(E_i)$ , we would have obtained the refined result that F + E' has snc support.

Set  $\nu := \pi \circ \mu : X'' \to X$ , and write  $K_{X''} = \nu^*(K_X + D) + E_{\nu}$ . We know that  $E_{\nu}$  is a divisor with snc support and coefficients  $\geq -1$ , and by the construction above, there exists a snc divisor F on X'' lying above  $X_{\text{sing}} \cup \text{Supp}(D)$  such that  $F + (E_{\nu})_{\text{red}}$  has snc support and  $\nu^*(K_X + D) - F$  is ample. Applying Theorem 4.6, we get:

Corollary 4.7. — Let (X, D) be a log canonical pair such that  $K_X + D$  is ample. Then the Kähler-Einstein metric  $\omega_{KE}$  on (X, D) is smooth on  $X_{reg} \setminus Supp(D)$ .

As we shall see in the course of the proof (cf §4.5.2), we do not obtain very precise estimates on the potential of the solution, even at order zero. However, it is tempting to believe that the potential  $\phi_{KE}$  of the Kähler-Einstein metric should be locally bounded outside of the non-klt locus of (X, D) defined as the support of the sheaf  $\mathcal{O}_X/\mathscr{I}(X, D)$  where  $\mathscr{I}(X, D)$  is the multiplier ideal sheaf of (X, D) (cf. e.g. [Laz04]). However, as this locus cannot be read easily on some log resolution, it does not seem obvious how one should tackle this question.

**4.4. Preliminaries: the regularized equation.** — We now borrow the notations of Theorem 4.6, and we let  $\omega_0$  be a Kähler form on X; it will be our reference metric in the following. Recall that we want to solve the equation

$$MA(\varphi) = \frac{e^{\varphi}dV}{\prod_{i} |s_{i}|^{2a_{i}}}$$

where the unknown function is  $\varphi$  a  $\theta$ -psh function,  $s_i$  are non-zero sections of  $\mathcal{O}_X(D_i)$ ,  $|\cdot|_i$  are smooth hermitian metrics on  $\mathcal{O}_X(D_i)$ ,  $f \in \mathscr{C}^{\infty}(X)$  and dV is a smooth volume form on X. Moreover, the expression  $\mathrm{MA}(\varphi)$  has to be understood as the non-pluripolar Mongeampère operator. It will be convenient for the following to differentiate the "klt part" of D from its "lc part", so we introduce the following notation:

$$D = \underbrace{\sum_{a_j < 1} a_j D_j}_{D_{klt}} + \underbrace{\sum_{a_k = 1} D_k}_{D_{lc}}$$

By Theorem 3.4, we know that the solution is the limit of any sequence of solutions of some appropriate regularized equations. The regularization process we are going to use concerns both the *a priori* non-Kähler class  $\{\theta\}$  and the "klt part" in the volume form:  $\prod_{a_i \le 1} |s_i|^{-2a_j}$ . More concretely, we will be studying the following equation:

(4.1) 
$$\langle (\theta + t\omega_0 + dd^c \varphi_{t,\varepsilon})^n \rangle = \frac{e^{\varphi_{t,\varepsilon} + f} dV}{\prod_{a_i < 1} (|s_i|^2 + \varepsilon^2)^{a_i} \prod_{a_k = 1} |s_k|^2}$$

Smoothness of the regularized solution. — At this point, it is still not completely clear that the solution  $\varphi_{t,\varepsilon}$  of equation (4.1) is smooth on  $X \setminus D_{lc}$ . So we translate our equation into the logarithmic setting: we set

$$\omega_{t,lc} := \theta + t\omega_0 - \sum_{a_k = 1} dd^c \log(\log|s_k|^2)^2$$

We may choose the hermitian metrics  $|\cdot|_k$  such that  $|s_k| < 1$  and such that  $\omega_{t,lc}$  defines a Kähler metric on  $X \setminus D_{lc}$  (cf [CG72, Gri76] e.g.). Of course this rescaling will depend on t, but we will explain how to bypass this problem later.

So using this new reference metric, one may rewrite equation (4.1) in the following form:

$$(\omega_{t,lc} + dd^c \psi_{t,\varepsilon})^n = \frac{e^{\psi_{t,\varepsilon} + f_t} \omega_{t,lc}^n}{\prod_{a_j < 1} (|s_j|^2 + \varepsilon^2)^{a_j}}$$

where  $\psi_{t,\varepsilon} = \varphi_{t,\varepsilon} + \sum_{a_k=1} \log(\log|s_k|^2)^2$  and  $f_t = -\log\left(\frac{\prod_k |s_k|^2 \log^2|s_k|^2 \omega_{t,lc}^n}{dV}\right)$ . Clearly,  $f_t$  is bounded (but only the lower bound is uniform in t) and smooth on  $X \setminus D_{lc}$ , but we know by Lemma 4.3 that  $f_t$  is smooth when read in the quasi-coordinates adapted to the pair  $(X, D_{lc})$ . Therefore, using the Theorem of Kobayashi and Tian-Yau (see Theorem 4.5), we know that the solution  $\psi_{t,\varepsilon}$  is bounded on  $X \setminus D_{lc}$ : there exists  $C_{t,\varepsilon} > 0$  such that

$$(4.2) -C_{t,\varepsilon} - \sum_{a_k=1} \log(\log|s_k|^2)^2 \leqslant \varphi_{t,\varepsilon} \leqslant C_{t,\varepsilon} - \sum_{a_k=1} \log(\log|s_k|^2)^2$$

Moreover,  $\psi_{t,\varepsilon}$  is smooth in the quasi-coordinates. In particular,  $\omega_{t,lc} + dd^c \psi_{t,\varepsilon}$  is a Kähler metric with bounded geometry on  $X \setminus D_{lc}$  and with Poincaré type growth along  $D_{lc}$ . Therefore it is complete and has a bounded curvature tensor. To prove the regularity theorem, we will thus have to obtain on each compact subset of  $X_0$  estimates on the potential  $\varphi_{\varepsilon}$  at

any order.

A first attempt at the uniform estimate. — The previous observation allows us to apply Yau's maximum principle (cf Proposition 4.4), and obtain that

$$\sup_{X \setminus D_{lc}} \psi_{t,\varepsilon} \leqslant \sup_{X \setminus D_{lc}} \left( \sum a_i \log(|s_i|^2 + \varepsilon^2) - f_t \right)$$

and similarly  $\inf_{X\setminus D_{lc}} \psi_{t,\varepsilon} \leqslant \inf_{X\setminus D_{lc}} \left(\sum a_i \log(|s_i|^2 + \varepsilon^2) - f_t\right)$ . If some coefficient  $a_i$  is negative, then we cannot obtain a bound for  $\sup \psi_{t,\varepsilon}$ . As for the lower bound,  $-f_t$  is not uniformly bounded from below because  $\omega_{t,lc}$  degenerates at t=0 (and if some  $a_i$  is positive,  $a_i \log(|s_i|^2 + \varepsilon^2)$  is not uniformly bounded below neither), so we cannot expect to find a lower bound for  $\psi_{t,\varepsilon}$  using this strategy. Therefore we need another method to obtain a zero-order estimate on the potential of the solution. In fact, we will need to add some barrier function to gain positivity, in the spirit of Tsuji's trick [Tsu88b] for the Laplacian estimate of a degenerate class; the novelty in our situation is that this is also needed for the zero-order estimates (as opposed to the klt case).

# **4.5.** Uniform estimate. — Before going any further, let us fix some notations.

4.5.1. A new framework. — We will denote by  $s_i$ ,  $i \in I$  (non-zero) sections of the (reduced) components of the divisor  $D_{red} + E$ , and by  $s_{\alpha}$ ,  $\alpha \in A$  (non-zero) sections of the (reduced) components of E; we endow all these line bundles with suitable smooth hermitian metrics (se below). Finally, we set  $X_0 := X \setminus (\operatorname{Supp}(D) \cup \operatorname{Supp}(E))$ , and define  $F := (D_{red} + E)_{red}$  as the reduced divisor  $X \setminus X_0$ .

The idea is to work on  $X_0$ . Of course, if we endow the last space with the Kähler metric  $\omega_{t,lc}$ , it will not be complete (near  $D_{klt}$  e.g.), so we won't be able to use Yau's maximum principle. Instead, we will rather use the following metric:

$$\omega_{\chi} := \theta + t\omega_0 - \sum_{i \in I} dd^c \log(\log|s_i|^2)^2 + dd^c \chi$$

where  $\chi := \sum_{\alpha} c_{\alpha} \log |s_{\alpha}|^2$  (recall that the  $c_{\alpha}$ 's are the coefficients of E).

We do here a slight abuse of notation because  $\omega_{\chi}$  depends on t. However, the following lemma shows that the dependence is harmless:

**Lemma 4.8.** — Up to changing the previously chosen hermitian metrics, the (1,1)-form  $\omega_{\chi}$  defines on  $X_0$  a smooth Kähler metric with Poincaré growth along F having bounded geometry, all of those properties being satisfied uniformly in t.

What we mean by this statement is that there exists a Poincaré-type metric  $\omega_P$  on  $X_0$  and a constant C>0 independent of t such that  $C^{-1}\omega_P\leqslant\omega_\chi\leqslant C\omega_P$ , and that in the appropriate quasi-coordinates attached to the pair (X,F), the coefficients  $g_{i\bar{j}}$  of  $\omega_\chi$  satisfy  $\left|\frac{\partial^{|\alpha|+|\beta|}g_{i\bar{j}}}{\partial z^\alpha\partial\bar{z}^\beta}\right|\leqslant C_{\alpha,\beta}$  for some constants  $C_{\alpha,\beta}>0$  independant of t. In particular,  $\omega_\chi$  has a uniformly (in t) bounded curvature tensor.

Proof of Lemma 4.8. — We know that  $\{\theta\} - c_1(E)$  is ample. Therefore, up to changing the hermitian metrics  $h_{\alpha}$  on  $\mathcal{O}_X(E_{\alpha})$ , we may suppose that  $\eta := \theta - \sum c_{\alpha}\Theta_{h_{\alpha}}(E_{\alpha})$  defines a smooth Kähler form on X (we designated by  $\Theta_{h_{\alpha}}(E_{\alpha})$  the curvature form of the hermitian line bundle  $(E_{\alpha}, h_{\alpha})$ ). Therefore, on  $X_0$ , we have:

$$\omega_{\chi} = \eta + t\omega_0 - \sum_{i \in I} dd^c \log(\log|s_i|^2)^2$$

and the statement follows easily from the computations of [CG72, Proposition 2.1] and [Kob84, Lemma 2] or [TY87].

4.5.2. Getting the lower bound. — First of all, we will use the crucial information that  $\varphi_{t,\varepsilon}$  converges (in the weak sense of distributions) to some  $\theta$ -psh function (cf first paragraph). By the elementary properties of psh (or quasi-psh) functions, we know that  $(\varphi_{t,\varepsilon})$  is uniformly bounded above on the compact set X (see e.g. [Hör94, Theorem 3.2.13]). Therefore, we obtain some uniform constant C such that

Now, recall that we chose  $\omega_{\chi}$  to be the new reference metric, so our equation becomes

$$(4.4) \qquad (\omega_{\chi} + dd^{c} u_{t,\varepsilon})^{n} = e^{u_{t,\varepsilon} + G_{\varepsilon}} \omega_{\chi}^{n}$$

where

$$u_{t,\varepsilon} := \varphi_{t,\varepsilon} + \sum_{i \in I} \log(\log|s_i|^2)^2 - \chi$$

and

$$G_{\varepsilon} := \chi + f + \sum_{a_i < 1} \log \left( \frac{|s_j|^2}{(|s_j|^2 + \varepsilon^2)^{a_j}} \right) - \log \left( \frac{\prod_{i \in I} |s_i|^2 \log^2 |s_i|^2 \omega_{\chi}^n}{dV} \right)$$

Here again we should mention that  $G_{\varepsilon}$  also depends on t through the last term involving  $\omega_{\chi}$ . For our purpose, we can ignore this dependence in order to simplify the notations.

We can see from (4.3) and Lemma 4.8 that  $G_{\varepsilon}$  has a uniform (in t and  $\varepsilon$ ) upper bound on  $X_0$ :

$$\sup_{\mathbf{v}} G_{\varepsilon} \leqslant C$$

Moreover, we know from (4.2) that  $\varphi_{t,\varepsilon} + \sum_{a_k=1} \log(\log|s_k|^2)^2$  is bounded. Therefore, it follows immediately that  $u_{t,\varepsilon}$  is bounded from below (but a priori non uniformly). Applying Yau's maximum principle (cf. Proposition 4.4) to the smooth function  $u_{t,\varepsilon}$  on the complete Kähler manifold  $(X_0, \omega_\chi)$  ensures that  $\inf_{X_0} u_{t,\varepsilon} \ge -\sup_{X_0} G_\varepsilon \ge -C$ . In terms of  $\varphi_{t,\varepsilon}$ , and recalling inequality (4.3) we get:

$$C \geqslant \varphi_{t,\varepsilon} \geqslant \chi - C - \sum_{i \in I} \log(\log|s_i|^2)^2$$

**4.6.** Laplacian estimate. — For the Laplacian estimate, we still work on the open manifold  $X_0$ . We endow it with the complete Kähler metric  $\omega_{\chi}$ , and we recall from Lemma 4.8 that  $\omega_{\chi}$  has uniformly bounded (bisectional) curvature.

As usual when one wants to compare to Kähler metrics  $\omega$  and  $\omega'$ , the strategy is to use an inequality of the form  $\Delta F \geqslant G$ , where F,G involve terms like  $\operatorname{tr}_{\omega}\omega'$ ,  $\operatorname{tr}_{\omega'}\omega$  or the local potentials of  $\omega' - \omega$ . There exist several variants of such inequalities, due e.g. to Chen-Lu, Yau, Siu, etc. involving different assumptions on the curvature of the metrics involved. In our case, as we have a control on the bisectional curvature of the reference metric  $\omega_{\chi}$ , on the Ricci curvature of the "unknown metric"  $\omega_{\chi} + dd^c u_{t,\varepsilon}$ , and on the laplacian  $\Delta_{\omega_{\chi}} G_{\varepsilon}$ , we could use any of these formulas.

We have chosen to use a variant of Siu's inequality [Siu87, p.99], which can be found in [CGP11, Proposition 2.1] (see also [Pău08, BBE+11]); notice the important feature allowing the factor  $e^{-F_{-}}$  for  $F_{-}$  quasi-psh which is crucial for us since the RHS of our Monge-Ampère equation has zeroes:

**Proposition 4.9.** Let X be a Kähler manifold of dimension n,  $\omega$ ,  $\omega'$  two cohomologous Kähler metrics on X. We assume that  $\omega' = \omega + dd^c u$  with  $\omega'^n = e^{u+F_+-F_-}\omega^n$ . and that we have a constant C > 0 satisfying:

- (i)  $dd^c F_{\pm} \geqslant -C\omega$ ,
- (ii)  $\Theta_{\omega}(T_X) \geqslant -C\omega \otimes \operatorname{Id}_{T_X}$ .

Then there exist some constant A > 0 depending only on n and C such that

$$\Delta_{\omega'}(\log \operatorname{tr}_{\omega}\omega' - Au + F_{-}) \geqslant \operatorname{tr}_{\omega'}\omega - nA.$$

Moreover, if one assumes that  $\sup F_+ \leqslant C, u \geqslant -C$  and that  $\log \operatorname{tr}_{\omega} \omega' - Au + F_-$  attains its maximum on X, then there exists M > 0 depending on n and C only such that:

$$\omega' \leq Me^{Au-F_-}\omega$$
.

Here  $\Delta$  (resp.  $\Delta'$ ) is the laplacian with respect to  $\omega$  (resp.  $\omega'$ ), and  $\Theta_{\omega}(T_X)$  is the Chern curvature tensor of the hermitian holomorphic vector bundle  $(T_X, \omega)$  (which may be identified with the tensor of holomorphic bisectional curvatures, usually denoted by the letter R).

Sketch of proof. — Siu's inequality applied to  $\omega = \sum g_{i\bar{j}} dz_i \wedge d\bar{z}_k$  and  $\omega' = \sum g'_{i\bar{j}} dz_i \wedge d\bar{z}_k$  yields:

$$\Delta'(\log \operatorname{tr}_{\omega}\omega') \geqslant \frac{1}{\operatorname{tr}_{\omega}\omega'}(-g^{j\bar{i}}R_{j\bar{i}} + \Delta(u + F_{+} - F_{-}) + g'^{k\bar{l}}R_{k\bar{l}}^{j\bar{i}}g'_{j\bar{i}})$$

Recollecting terms coming (with different signs) from the scalar and the Ricci curvature, we will obtain a similar inequality involving only a lower bound for the holomorphic bisectional curvature, namely

(4.5) 
$$\Delta' \log \operatorname{tr}_{\omega} \omega' \geqslant \frac{\Delta(u + F_{+} - F_{-})}{\operatorname{tr}_{\omega}(\omega')} - B \operatorname{tr}_{\omega'} \omega$$

where B is a lower bound for the bisectional curvature of  $\omega$ : this is the content of [CGP11, Lemma 2.2].

Clearly,  $\Delta u = \operatorname{tr}_{\omega} \omega' - n$  so that  $\Delta(u + F_+) \geqslant -n(C + 1)$ . As  $\operatorname{tr}_{\omega} \omega' \operatorname{tr}_{\omega'} \omega \geqslant n$ , we get

(4.6) 
$$\frac{\Delta(u+F_{+})}{\operatorname{tr}_{\omega}\omega'} \geqslant -(1+C)\operatorname{tr}_{\omega'}\omega$$

As for the second laplacian, we write

$$0 \leqslant C\omega + dd^c F_- \leqslant \operatorname{tr}_{\omega'}(C\omega + dd^c F_-)\omega'$$

and we take the trace with respect to  $\omega$ :

$$\frac{nC + \Delta_{\omega} F_{-}}{\operatorname{tr}_{\omega} \omega'} \leqslant C \operatorname{tr}_{\omega'} \omega + \Delta_{\omega'} F_{-}$$

so that

(4.7) 
$$\Delta_{\omega'} F_{-} \geqslant \frac{\Delta_{\omega} F_{-}}{\operatorname{tr}_{\omega'} \omega'} - C \operatorname{tr}_{\omega'} \omega$$

Plugging (4.6) and (4.7) into (4.5), we get:

$$\Delta'(\log \operatorname{tr}_{\omega}\omega' + F_{-}) \geqslant -C_1\operatorname{tr}_{\omega'}\omega$$

for  $C_1 = 1 + B + 2C$ . Finally, using  $\Delta' u = n - \operatorname{tr}_{\omega'} \omega$ , we see that

$$\Delta'(\log \operatorname{tr}_{\omega}\omega' - (C_1 + 1)u + F_-) \geqslant \operatorname{tr}_{\omega'}\omega - n(C_1 + 1)$$

which shows the first assertion by chosing  $A := 1 + C_1$ .

As for the second part, if we denote by p the point where the maximum is attained, then one has  $(\operatorname{tr}_{\omega'}\omega)(p) \leqslant C_2$ . Using the basic inequality  $\operatorname{tr}_{\omega}\omega' \leqslant e^{u+F_+-F_-}(\operatorname{tr}_{\omega'}\omega)^{n-1}$ , one gets

$$\log(\operatorname{tr}_{\omega}\omega') = (\log \operatorname{tr}_{\omega}\omega' - Au + F_{-}) + Au - F_{-}$$

$$\leq (u(p) + F_{+}(p) - F_{-}(p)) + (n-1)\log(nA) - Au(p) + F_{-}(p) + Au - F_{-}$$

$$\leq C_{2} + Au - F_{-}$$

where  $C_2 = \sup F_+ + (n-1)\log(nA) - (A-1)\inf u$  (recall that A can be chosen to be positive). This concludes the proof of the proposition.

Recall that we are interested in equation (4.4) given by

$$(\omega_{\chi} + dd^{c}u_{t,\varepsilon})^{n} = e^{u_{t,\varepsilon} + G_{\varepsilon}}\omega_{\chi}^{n}$$

We obtained the zero-order estimate on  $u_{t,\varepsilon}$  in the last section, and now we want a Laplacian estimate. In order too use the previous proposition we first have to decompose  $G_{\varepsilon}$  as a difference of  $C\omega_{\chi}$ -psh functions in order to use the result above. Recall that

$$G_{\varepsilon} := \chi + f + \sum_{a_i < 1} \log \left( \frac{|s_j|^2}{(|s_j|^2 + \varepsilon^2)^{a_j}} \right) - \log \left( \frac{\prod_{i \in I} |s_i|^2 \log^2 |s_i|^2 \omega_{\chi}^n}{dV} \right)$$

By [Kob84] or [Gue12b, Lemma 1.6], the last term is already known to be smooth in the quasi-coordinates (and it depends neither on t nor on  $\varepsilon$ ). We claim that

$$G_{\varepsilon} = \underbrace{\left[\chi + f + \sum_{a_j < 1} \log|s_j|^2 + \sum_{a_j < 0} \log(|s_j|^2 + \varepsilon^2)^{-a_j}\right]}_{F_+} - \underbrace{\left[\sum_{0 < a_j < 1} \log(|s_j|^2 + \varepsilon^2)^{a_j}\right]}_{F_-}$$

gives the desired decomposition  $G_{\varepsilon} = F_{+} - F_{-}$  in the notations of Proposition 4.9. Indeed,  $\chi, f, \log |s_{i}|^{2}$  are quasi-psh, thus  $C\omega_{\chi}$ -psh for some uniform C > 0 as  $\omega_{\chi}$  dominates some fixed Kähler form, cf Lemma 4.8. Moreover, a simple computation leads to the identity:

$$dd^c \log(|s|^2 + \varepsilon^2) = \frac{\varepsilon^2}{(|s|^2 + \varepsilon^2)^2} \cdot \langle D's, D's \rangle - \frac{|s|^2}{|s|^2 + \varepsilon^2} \cdot \Theta_h$$

where  $\Theta_h$  is the curvature of the hermitian metric h implicit in the term  $|s|^2$ , and D's is the (1,0)-part of Ds where D is the Chern connection attached to  $(\mathcal{O}_X(\operatorname{div}(s)),h)$ . If follows that  $F_{\pm}$  are  $C\omega_X$ -psh for some uniform C>0.

We can now apply Proposition 4.9 to the setting:  $\omega = \omega_{\chi}, \omega' = \omega_{\chi} + dd^{c}u_{t,\varepsilon}, F_{+} - F_{-} = G_{\varepsilon}$ . Indeed, it is clear that  $F_{+}$  is uniformly upper bounded, we just saw that  $F_{\pm}$  are  $C\omega_{\chi}$ -psh, and we know from the previous section that  $u_{t,\varepsilon}$  has a uniform lower bound. Furthermore,  $\log \operatorname{tr}_{\omega}\omega' - Au + F_{-}$  attains its maximum on  $X_{0}$ : indeed, -Au tends to  $-\infty$  near the boundary of  $X_{0}$ ,  $F_{-}$  is bounded (it is even smooth), and  $\operatorname{tr}_{\omega}\omega' = \Delta_{\omega_{\chi}}(\varphi_{t,\varepsilon} + \sum_{i \in I} \log(\log |s_{i}|^{2})^{2} - \chi)$  is bounded on  $X_{0}$  (we know it for the term  $\Delta_{\omega_{\chi}}(\varphi_{t,\varepsilon} - \chi)$  and it is an elementary computation for the other term).

In conclusion, we may use Proposition 4.9 to obtain the following estimate:

$$\theta + t\omega_0 + dd^c \varphi_{t,\varepsilon} \leqslant M \left( \prod_{i \in I} \log^2 |s_i|^2 \right)^C \cdot \prod_{\alpha \in A} |s_\alpha|^{-c_\alpha \cdot C} \prod_{a_j > 0} |s_j|^{-2a_j} \, \omega_\chi$$

For the "reverse inequality", we use the identity

$$(\omega_{\chi} + dd^{c}u_{t,\varepsilon})^{n} = e^{u_{t,\varepsilon} + G_{\varepsilon}}\omega_{\chi}^{n}$$

which leads to the inequality

$$\operatorname{tr}_{\omega_{\chi}+dd^{c}u_{t,\varepsilon}}\omega_{\chi} \leqslant e^{-(u_{t,\varepsilon}+G_{\varepsilon})} \left(\operatorname{tr}_{\omega_{\chi}}(\omega_{\chi}+dd^{c}u_{t,\varepsilon})\right)^{n-1}$$

and therefore

$$\theta + t\omega_0 + dd^c \varphi_{t,\varepsilon} \geqslant M^{-1} \prod_{a_j < 1} \frac{|s_j|^2}{(|s_j|^2 + \varepsilon^2)^{a_j}} \cdot \left( \prod_{i \in I} \log^2 |s_i|^2 \right)^{-C} \cdot \prod_{\alpha \in A} |s_\alpha|^{c_\alpha \cdot C} \prod_{a_j > 0} |s_j|^{2a_j} \omega_\chi$$

for some uniform C, M > 0 (different from the previous ones).

In particular, for any compact set  $K \in X_0$ , there exists a constant  $C_K > 0$  satisfying

$$C_K^{-1} \omega_0 \leqslant \theta + t\omega_0 + dd^c \varphi_{t,\varepsilon} \leqslant C_K \omega_0$$

Using Evans-Krylov theorem and the classical elliptic theory shows that the potential  $\varphi_{t,\varepsilon}$  satisfies uniform  $\mathscr{C}^{k,\alpha}$  estimates on any  $\Omega \subseteq K$  for each  $k,\alpha$ . Thus the theorem is proved.

**Remark 4.10.** — One can easily obtain somewhat more precise estimates. Indeed, if  $\alpha$  is a nef and big class and E an effective  $\mathbb{R}$ -divisor such that  $\alpha - E$  is ample, then we have in fact that for every  $\delta > 0$ ,  $\alpha - \delta E$  is ample (write  $\alpha - \delta E = (1 - \delta)\alpha + \delta(\alpha - E)$ ). Applying this observation to Theorem 4.6, we see that for every  $\delta > 0$ , there exists  $C_{\delta} > 0$  such that:

$$\varphi_{\text{KE}} \geqslant \delta \sum_{\alpha} c_{\alpha} \log |s_{\alpha}|^2 - \sum_{i \in I} \log(\log |s_i|^2)^2 - C_{\delta}$$

for every  $\delta > 0$ . One could apply the same argument to the Laplacian estimates.

**4.7.** The nef and big case. — Let us now assume that (X, D) is a log canonical pair satisfying that  $K_X + D$  is merely nef and big. We claim that all the previously obtained results remain true *mutatis mutandis*. More precisely, here is the theorem:

**Theorem 4.11.** — Let X be a compact Kähler manifold of dimension n, dV some volume form,  $D = \sum a_i D_i$  a  $\mathbb{R}$ -divisor with coefficients in  $(-\infty, 1]$  and defining sections  $s_i$ ,  $E = \sum c_{\alpha} E_{\alpha}$  an effective  $\mathbb{R}$ -divisor such that  $D_{red} + E$  has snc support, and  $\beta \in H^{1,1}(X, \mathbb{R})$  a nef and big cohomology class such that  $\beta - c_1(E)$  is ample. Let also  $\theta \in \beta$  be a smooth representative. Then the  $\theta$ -psh function  $\varphi$  with full Monge-Ampère mass solution of

$$\langle (\theta + dd^c \phi)^n \rangle = \frac{e^{\varphi} dV}{\prod_i |s_i|^{2a_j}}$$

is smooth outside of  $Supp(D) \cup Supp(E)$ .

Proof. — The existence is guaranteed by the results in the second section. Choose now  $\theta \in \beta$  a smooth representative, and  $\omega_0$  a Kähler form on X. We know that for each t > 0, there exists a smooth function  $f_t$  on X such that  $\omega_t := \theta + t\omega_0 + dd^c f_t$  is a Kähler form on X. Therefore, one can produce a metric of Carlson-Griffiths type  $\omega_{t,lc} = \omega_t + \sum_{a_k=1} \log(\log|s_k|^2)^2$ , and solve the "half-regularized" equation as in the semipositive case. So we managed to write the solution as limit of Kähler metric on  $X \setminus D_{lc}$  being smooth in the quasi-coordinates. The end of the proof is exactly the same as in the semipositive case.

As for the ample case, the discussion about finding an appropriate log resolution  $\pi$  of (X, D) where the augmented base locus of  $\pi^*(K_X + D)$  is a snc divisor lying above  $X_{\text{sing}} \cup \text{Supp}(D) \cup \mathbb{B}_+(K_X + D)$  applies, so that we have:

Corollary 4.12. — Let (X, D) be a log canonical pair such that  $K_X + D$  is nef and big. Then the Kähler-Einstein metric  $\omega_{KE}$  on (X, D) is smooth on  $(X_{reg} \setminus Supp(D)) \cap Amp(K_X + D)$ . As we already mentioned it in the introduction, the base point free theorem is false in the general setting of log canonical pairs; having  $K_X + D$  nef and big does not imply that it is semi-ample. Therefore the previous statement is more general than Theorem 4.6. This being said, the remaining challenging case is the one of a log canonical pair (X, D) of log general type, ie satisfying  $K_X + D$  big.

If the pair is klt, then its admits a log canonical model by the seminal results of [BCHM10], where the finiteness of the (log) canonical ring was established. The smoothness, on the appropriate locus (ie  $X_0 := X_{\text{reg}} \setminus (\text{Supp}(D) \cup \mathbb{B}_+(K_X + D))$ ), of the Kähler-Einstein metric constructed in [EGZ09] follows from the existence of the log canonical model, as one can then reduce the problem to the semipositive case (cf also [BEGZ10] or [Gue12a] for discussions about this topic). But in the general log canonical case, the existence of a log canonical model – which would, by the same reduction argument combined with our regularity results in the semi-ample case, imply the smoothness of the Kähler-Einstein metric on  $X_0$  – is still a major conjecture. Accordingly, we do not know if the general (singular) Kähler-Einstein metric that we constructed using variational methods is smooth on  $X_0$ . Even in the klt case, showing the smoothness of the KE metric by direct analytical methods remains a challenging problem.

**4.8.** Completeness. — Given a log canonical pair (X, D) such that  $K_X + D$  is ample, we know that there exists a unique Kähler-Einstein current in  $c_1(K_X + D)$  inducing a genuine Kähler-Einstein metric  $\omega_{\text{KE}}$  on  $X_0 = X_{\text{reg}} \setminus \text{Supp}(D)$  with maximal volume. In this section, we will try to understand whether the Kähler manifold  $(X_0, \omega_{\text{KE}})$  is complete. Let us first begin with the already known cases.

4.8.1. Previous results. — We choose a log resolution  $\pi: X' \to X$  of (X, D) and write  $K_{X'} = \pi^*(K_X + D) - E$ , where  $E = \sum a_i E_i$  is a divisor with snc support and coefficients  $a_i \leq 1$ .

Let us begin by the easiest case of a log smooth pair; ie  $\pi = \text{Id}_X$ . Then, if E = D is reduced, we known from [Kob84, TY87] that  $\omega_{\text{KE}}$  has Poincaré singularities along E and hence is complete.

In fact, if (X, D) is a purely log canonical pair in the sense that for some resolution, E is reduced, then Yau showed in [Yau93] that the Kähler-Einstein metric is complete.

Coming back to the log smooth case, then if  $E = \sum (1 - \frac{1}{m_i})E_i$  has orbifold coefficients (ie  $m_i \in \mathbb{Z}_{>0} \cup \{+\infty\}$ ), it follows from [**TY87**] that  $\omega_{\text{KE}}$  is smooth in the orbifold sense along the strictly orbifold part of E (ie the  $E_i$  such that  $m_i < +\infty$ ), and therefore  $\omega_{\text{KE}}$  is complete if and only if E is reduced.

If now  $E = \sum a_i E_i$  with  $a_i \geqslant 1/2$  then it follows from [Gue12b] that  $\omega_{\text{KE}}$  has mixed cone and Poincaré singularities; as cone metrics are never complete, one observes again that  $\omega_{\text{KE}}$  is complete if and only if E is reduced.

Before we get to the main result of this section, let us give the two main tools we are going to use. The first one is the following consequence of Yau-Schwarz lemma in its version for volume forms [Yau78a, Theorem 3]:

**Theorem 4.13.** — Let X be a complex manifold of dimension n endowed with two Kähler metrics  $\omega_1$  and  $\omega_2$ . One assumes that the Ricci curvature of  $\omega_1$  is bounded from below by some constant  $-K_1$ , and that  $\omega_2$  has scalar curvature bounded from above by some negative constant  $K_2$ , and Ricci curvature bounded from below in a non-quantitative way. If  $\omega_1$  is complete, then we have:

$$\omega_2^n \leqslant \frac{K_1}{K_2} \, \omega_1^n$$

The second tool is the following consequence of the fundamental results of [BCHM10] (cf e.g. [EGZ09, BEGZ10, Gue12a]):

**Theorem 4.14.** — If (X, D) is a klt pair of log general type (ie  $K_X + D$  is big), then the Kähler-Einstein metric of (X, D) is smooth on  $X_{reg} \setminus (\operatorname{Supp}(D) \cup \mathbb{B}_+(K_X + D))$ .

4.8.2. Completeness of the Kähler-Einstein metric. — The goal of this section is to prove the following:

**Theorem 4.15**. — Let (X, D) be a log canonical pair such that  $K_X + D$  is ample, and let  $\omega_{KE}$  be the smooth Kähler-Einstein metric on  $X_0$ . Then:

- (i) If (X, D) is log smooth, then  $\omega_{KE}$  is complete if and only if D is reduced.
- (ii) If (X, D) is klt, then  $\omega_{KE}$  is not complete, unless X is smooth and D = 0.

In fact, we will see that point (ii) is valid for any klt pair (X, D) with  $K_X + D$  big.

**Remark 4.16.** — In the positively curved case, Bonnet-Myers theorem shows that the Kähler-Einstein metrics are never complete (besides, the singularities are always klt by [BBE+11, Proposition 3.8]).

Proof of Theorem 4.15. — Let us begin with the first point, and assume that  $\omega_{\text{KE}}$  is complete for a log smooth pair (X, D). We write  $D = D_{klt} + D_{lc}$  as usual, and we suppose that  $D_{klt} \neq 0$ . Let us choose some small open subset U of  $X_{lc} = X \setminus D_{lc}$  where  $D_{klt} = aH$  is given by one smooth hypersurface H with section s defined on X. Then we choose  $\varepsilon > 0$  small enough, so that  $K_X + D + \varepsilon H$  is still ample. Then there exists a Kähler-Einstein metric  $\omega$  (with Ricci curvature equal to -1) for the pair  $(X, D + \varepsilon H)$ . We know from Lemma 4.17 that the local potentials  $\varphi_{\text{KE}}$  (resp.  $\varphi$ ) of  $\omega_{\text{KE}}$  (resp.  $\omega$ ) satisfy on U:

$$C_0 \geqslant \varphi_{\text{KE}}, \quad \varphi \geqslant -\log(\log^2|s|^2) - C_0$$

for some constant  $C_0 > 0$  depending only on U. Therefore, we have:

$$\omega^n \geqslant C_0 |s|^{-a-\varepsilon} \log^{-2} |s|^2$$

so that  $\omega^n \geqslant C_1|s|^{-a-\varepsilon/2}$  for some  $C_1 > 0$  depending only on U. By Yau-Schwarz lemma, we should have  $\omega^n \leqslant \omega_{\text{KE}}^n$ . But as  $\varphi_{\text{KE}}$  is bounded from above on U, we also have  $\omega_{\text{KE}}^n \leqslant C_2|s|^{-a}$  on U, which is a contradiction.

Let us get to point (ii). We may assume that (X, D) is log smooth with  $K_X + D$  semipositive and big. We choose a smooth hypersurface H which is a component of D. As for  $\varepsilon > 0$  small enough,  $(X, D + \varepsilon H)$  is klt and of log general type, we also get a Kähler-Einstein metric  $\omega$  with -1 Ricci curvature smooth (at least) on the same locus as  $\omega_{KE}$  (because  $\mathbb{B}_+(K_X + D + \varepsilon H) \subset \mathbb{B}_+(K_X + D) \cup H$ ). Its potentials are known to have minimal singularities [BEGZ10, Gue12a] and in particular they have zero Lelong numbers. We can then get a contradiction exactly as in the previous case using Yau-Schwarz lemma.

Let us now explain the proof of the following lemma that we used to show (i) in the previous theorem:

**Lemma 4.17.** — Let (X, D) be a log smooth pair with  $K_X + D$  ample. Then on any relatively compact subset  $U \subseteq X \setminus D_{lc}$ , the local potentials  $\varphi_{KE}$  of the Kähler-Einstein metric satisfy:

$$C \geqslant \varphi_{\text{KE}} \geqslant -\sum_{i} \log \log^{2}(|s_{i}|^{2}) - C$$

where C > 0 depends only on U and the sum is taken over the section  $s_i$  of the components of  $D_{klt}$  restricted to U.

*Proof.* — This follows from the uniform estimate given in §4.5.2. Indeed, in the log smooth case, we do not need to use the barrier function  $\chi$  because  $K_X + D$  is already ample.  $\square$ 

We should mention that a simpler proof of the incompleteness of the Kähler-Einstein metric attached to a klt pair (X, D) with  $K_X + D$  ample (or trivial) exists: indeed, in a log-resolution, this metric  $\omega_{KE}$  is known to satisfy an estimate like  $\omega_{KE} \leq C_{\varepsilon} |s_E|^{-2\varepsilon} \prod_i |s_i|^{-2a_i} \omega$  where  $\omega$  is smooth, and the relative canonical divisor of the resolution is  $\sum a_i [s_i = 0]$ , and E is an snc exceptional divisor,  $\varepsilon > 0$  being arbitrary. Therefore the manifold  $(X_0, \omega_{KE})$  is easily seen to have finite diameter and hence cannot be complete unless X is smooth and D = 0. Note that this proof does not work neither in the big case, nor in the case (i) of Theorem 4.15.

Of course, Theorem 4.15 misses a lot of interesting cases. We conjecture that for a lc pair (X, D) with  $K_X + D$  ample, the Kähler-Einstein metric  $\omega_{KE}$  on  $X_0$  is complete if and only if (X, D) has purely log canonical singularities in the sense that the klt locus should coincide with  $X_0 = X_{reg} \setminus \text{Supp}(D)$ ; where we recall that  $\text{KLT}(X, D) := X \setminus \text{Supp}(\mathcal{O}_X/\mathscr{I}(X, D))$ ,  $\mathscr{I}(X, D)$  being the multiplier ideal attached to (X, D), cf [Laz04].

About uniqueness of the Kähler-Einstein metrics. — First of all, in the case where X is smooth and  $K_X$  is ample, then uniqueness of the Kähler-Einstein metric constructed by Aubin and Yau is a straightforward consequence of the maximum principle. Generalizing this principle to some complete Kähler manifolds in [Yau75], Yau could prove that on a

Kähler manifold, there can be only one *complete* Kähler metric  $\omega$  satisfying  $\operatorname{Ric} \omega = -\omega$ . In particular, this result has been applied by Kobayashi and Tian-Yau to show the uniqueness of the Kähler-Einstein metric for a log smooth pair (X, D) satisfying  $K_X + D$  ample, cf [Kob84, TY87]. Using Yau-Schwarz lemma in a more subtle way (through the notion of almost complete metric), they also show uniqueness when  $K_X + D$  is only assumed nef, big and ample modulo D, which means that  $K_X + D$  intersects positively every curve not contained in D. For example, if (X, D) is a log resolution of some canonically polarized singular variety, these assumptions are not satisfied.

In our situation we proceed in a different manner: we first use the volume assumption (as a replacement for completeness) to show that the Kähler-Einstein metric, originally defined on the (log) regular locus, extends to define a positive current on X whose local potentials glue to define a solution with full Monge-Amère mass of a global Monge-Ampère equation to which we can apply the comparison principle to finally deduce the uniqueness.

## 5. Applications

## 5.1. Yau-Tian-Donaldson conjecture for canonically polarized singular varieties.

— Let us start with the following converse of Theorem A stated in the introduction.

**Proposition 5.1.** — Let X be projective variety satisfying the conditions  $G_1$  and  $S_2$ , and such that  $K_X$  is  $\mathbb{Q}$ -ample. If X admits a Kähler-Einstein metric  $\omega$  in the following sense:  $\omega$  is a K-E metric on  $X_{\text{reg}}$  and its total volume there is equal to  $K_X^n$ , then X has semi-log canonical singularities.

*Proof.* — We begin with the case where X is normal.

Let us first show that  $\phi := \log \omega^n$  on  $X_{\text{reg}}$  extends to an element in  $\mathcal{E}(X, K_X)$ . By the K-E equation  $dd^c\phi := -\text{Ric }\omega = \omega \geqslant 0$  and hence, since X is normal,  $\phi$  extends to a positively curved singular metric on  $K_X$  over all of X. Thus, writing  $\omega = dd^c\psi$  for some other such metric  $\psi$  the compactness of X forces the K-E equation  $\text{MA}(\phi) = e^{\phi}$  (up to shifting  $\phi$  by a constant) globally on X (since the non-pluripolar MA does not charge the singular locus of X). Moreover, by the volume assumption we have that  $\int_X \text{MA}(\phi) = K_X^n$  and hence  $\phi \in \mathcal{E}(X, K_X)$ , as desired.

Next, fix a resolution  $\pi: X' \to X$  and assume, to a get a contradiction, that  $p^*K_X = K_{X'} + D$  where D is a snc  $\mathbb{Q}$ -divisor such that  $D = D' + (1 + \delta)E$  for some  $\delta > 0$ , where D' is a  $\mathbb{Q}$ -divisor and E is a smooth irreducible divisor transversal to the support of D'. Since  $\phi$  has maximal MA-mass it follows, as shown in  $[\mathbf{BBE^{+}11}]$  using an Izumi type estimate, that  $\pi^*\phi$  has no Lelong numbers. In particular, it follows from the characterization of Lelong numbers that there exists a neigbourhood U of E such that  $\pi^*\phi \geqslant \frac{1}{2}\delta \log |s_E|^2 - C$  in local trivializations. Moreover, we may take U such that D does not intersect U. But then it

follows from the K-E equation that

$$K_X^n \geqslant C' \int_U e^{\pi^* \phi - (1+\delta) \log|s_E|^2} \geqslant C'' \int_U e^{-(1+\delta/2) \log|s_E|^2} = \infty,$$

which gives the desired contradiction.

We move on to the general case when X is only assumed to be  $G_1$  and  $S_2$ . As we observed in §2.3, the result of Proposition 2.6 holds actually in the general  $G_1$  and  $S_2$  case (we did not use at all that the singularities were slc); however we should be careful and work instead on a log-resolution of  $(X^{\nu}, C_{X^{\nu}})$  because the formula  $\nu^*K_X = K_{X^{\nu}} + C_{X^{\nu}}$  could not be meaningful anymore if  $C_{X^{\nu}}$  is not Cartier, cf §2.3 and the remarks following the identity (2.1). So the first conclusion is that the weight  $\phi := \log \omega^n$  on  $X_{\text{reg}}$  extends on the normalization  $X^{\nu}$  to a psh weight in  $\mathcal{E}(X^{\nu}, \nu^*K_X)$ . Then we take a log-resolution  $\pi: X' \to X^{\nu}$  of the pair  $(X^{\nu}, C_{X^{\nu}})$  where  $C_{X^{\nu}}$  is the conductor of the normalization; a priori, this is just an effective divisor, possibly non-reduced. We write (X', D') for the new pair that we obtain on X'. Then same arguments as earlier show that D' has coefficients less than or equal to 1, which amounts to saying that  $(X^{\nu}, C_{X^{\nu}})$  is log canonical, or equivalently that X has semi-log canonical singularities.

To relate this to K-stability we recall that Odaka [Oda08] has shown that, if X is K-semistable, then X has semi-log canonical singularities (recall that we assume that X is  $G_1$  and  $G_2$  and that  $G_3$  and that  $G_4$  is K-stable [Oda11]. Hence, combining our results with Odaka's results gives the following confirmation of the Yau-Tian-Donaldson conjecture for varieties being  $G_1$  and  $G_2$ , with  $G_3$  ample:

**Theorem 5.2.** — Let X be a  $G_1$  and  $S_2$  projective variety such that  $K_X$  is ample. Then X admits a Kähler-Einstein metric iff X is K-stable.

It would be interesting to have a direct analytical proof of the implication "Kähler-Einstein implies K-stable" as shown in [Ber12] the (log) Fano case (where K-stability has to be replaced by K-polystability in the presence of holomorphic vector fields).

**5.2.** Automorphism groups of canonically polarized varieties. — The existence and uniqueness of Kähler-Einstein metrics established in Theorem A allows us to give an analytical proof of the following result shown in [BHPS12] (where two proofs were given, one cohomological and one geometric)

**Proposition 5.3.** — Let X be a normal stable variety. Then X admits no non-trivial infinitisimal automorphisms.

*Proof.* — By general results on automorphism groups of normal varieties (see [BBE<sup>+</sup>11, Lemma 5.2] and references therein) it is equivalent to show that any holomorphic vector field V on  $X_{\text{reg}}$  vanishes identically. To prove this vanishing we first observe that, by normality, V is the infinitisimal generator of a complex one-parameter family of automorphism F of X

and in particular of  $X_{\text{reg}}$ . Fix a Kähler-Einstein current  $\omega$  on X. By the naturality of the KE-equation it follows that  $F^*\omega$  is also a KE-current and hence by uniqueness  $F^*\omega = \omega$  on  $X_{\text{reg}}$ . Let us denote by  $V_r$  and  $V_i$  the real and imagnary parts of V, which are infinitisimal generators of real one parameter families of automorphisms that we will denote by  $F_r$  and  $F_i$  respectively, which, by the previous argument, also preserve  $\omega$ . Next, note that any automorphism automatically lifts to the line bundle  $K_X$  over  $X_{\text{reg}}$  and thus it follows from general principles that the real part  $V_r$  of V is a Hamiltonian vector field, i.e.  $i_{V_r}\omega = dh$  for some smooth function h on  $X_{reg}$ . But then, by Cartan's formula, the Lie deriviative  $L_{V_i}\omega$  is given by  $d(i_{W_r}\omega) = dJi_{V_r}\omega = dJdh = dd^ch$ . Since the flow  $F_i$  defined by  $V_i$  also preserves  $\omega$  it thus follows that  $dd^ch = 0$ . But by normality it follows that h = 0 (indeed, by normality h is bounded and we can thus apply the maximum principle on a resolution). Hence  $i_{V_r}\omega = 0$  on  $X_{reg}$ , which forces  $V_r = 0$  on  $X_{reg}$ , since  $\omega$  is Kähler there and in particular pointwise non-degenerate. Finally, by the same argument  $V_i = 0$  (for example replacing V with JV) and hence V = 0 as desired.

In the case when X is smooth there is a simple cohomological proof of the previous proposition: by Serre duality  $H^0(X, TX)$  is isomorphic to  $H^{n-1}(X, -K_X)$ , which is trivial by Kodaira vanishing (since  $K_X$  is ample). In the case when X is log canonical a similar cohomological argument can be used [BHPS12], relying on the Bogomolov-Sommese vanishing result for log canonical singularities, established in [GKKP11, Theorem 7.2]. Indeed, if V does not vanish identically then contracting with V on  $X_{\text{reg}}$  maps  $K_X$  to a rank one reflexive sheaf in Hom  $(K_X, \Omega_X^{[n-1]})$ , where  $\Omega_X^{[n-1]}$  is the sheaf of reflexive (n-1)-forms on X and hence by the Bogomolov-Sommese vanishing result in [GKKP11] the Kodaira dimension of  $K_X$  is at most n-1, which contradict the ampleness of  $K_X$ .

## 6. Outlook

**6.1. Towards Miyaoka-Yau type inequalities.** — For simplicity we will only consider the case n=2 (but a similar discussion applies in the general case). We set  $E:=\Omega_X^1$ , the cotangent bundle of X. The classical case is when X is smooth with  $K_X$  ample, where the Miyaoka-Yau inequality says

$$c_1(E)^2 \leqslant 3c_2(E).$$

Let us briefly recall Yau's differential-geometric proof. We equip E with the Hermitian metric induced by  $\omega$  and denote by  $(E,\omega)$  the corresponding Hermitian vector bundle. Then, if  $\omega$  is Kähler-Einstein a direct local calculation gives the *point-wise* inequality

$$c_1(E,\omega)^2 \leqslant 3c_2(E,\omega)$$

formulated in terms of the Chern-Weil representatives  $c_i(E,\omega)$  of the corresponding Chern classes. Hence, integrating immediately gives the Miyaoka-Yau inequality. Repeating this argument in the singular case when X a stable surface and using Theorem A gives the following

**Proposition 6.1.** — The following inequality holds for a stable surface equipped with the canonical Kähler-Einstein metric  $\omega$  on its regular part:

$$c_1(K_X)^2 \leqslant 3 \int_{X_{\text{reg}}} c_2(E,\omega)$$

with equality iff  $\omega$  has constant holomorphic sectional curvature, i.e.  $(X_{reg}, \omega)$  is locally isometric to a ball.

*Proof.* — Since the point-wise inequality above still holds, by the KE-condition, we can simply integrate it over  $X_{reg}$  and use that, by Theorem A,  $c_1(K_X)^2 = \int_{X_{reg}} c_1(E,\omega)^2$ . The conditions for equality are well-known in the point-wise inequality.

Since  $\omega$  is canonically attached to X one could simply define the rhs appearing in the inequality above as the "analytical second Chern number"  $c_{2,an}(X)$  of X. However, it should be stressed that it is not even a priori clear that  $c_{2,an}(X)$  is finite, even though we expect that this is the case. More precisely, we expect that  $c_{2,an}(X)$  can be identified with (or at least bounded from above) by a suitable algebraically defined second Chern class number  $c_2(X)$ . Various definitions of such Chern numbers have been proposed in the litterature and we refer the reader to the paper of Langer [Lan00] where very general algebraic Miyaoka-Yau type inequalities are obtained, which in particular apply to stable surfaces. More generally, as before, our arguments apply to log canonical pairs.

6.2. The Weil-Peterson geometry of the moduli space of stable varieties. — In this section we will briefly explain how the finite energy property of the Kähler-Einstien metric on a stable variety naturally appears in the geometric study of the moduli space  $\mathcal{M}$  of all stable varieties. In a nut shell, the Kähler-Einstein metrics on stable varieties induces a metric on the  $\mathbb{Q}$ -line bundle  $\mathcal{L} \to \mathcal{M}$  over the moduli space defined by the top Deligne pairing of  $K_X$  and the finite energy condition is precisely the condition which makes sure that the metric is point-wise finite. The relation to Weil-Petterson geometry comes from the well-known fact that the curvature form of the corresponding metric over the moduli space  $\mathcal{M}_0$  of all smooth stable varieties (i.e. all canonically polarized n-dimensional manifolds, with a fixed oriented smooth structure) coincides with the Weil-Peterson metric  $\Omega_{WP}$  on  $\mathcal{M}_0$  [FS90, Sch12].

To be a bit more precise we first recall that given a line bundle  $L \to X$  over a (complex) n-dimensional algebraic variety X its top Deligne pairing i.e. the (n+1)-fold Deligne pairing of L with itself is a complex line that we will denote by  $\langle L \rangle$  [Elk89, Elk90]. Equipping  $\langle L \rangle$  with an Hermitian metric  $\phi$  (using additive notation as before) induces a Hermitian metric  $\langle \phi \rangle$  on  $\langle L \rangle$ , satisfying the change of metric formula:  $\langle \phi \rangle - \langle \psi \rangle = (\mathcal{E}(\phi) - \mathcal{E}(\psi))$  (up to a multiplicative normalization constant), where  $\mathcal{E}$  is the energy functional appearing in §3 (compare [PS04]). Fixing a smooth reference metric  $\phi_0$  on L one can use the latter transformation formula to define the metric  $\langle \phi \rangle$  as long as  $\phi$  has finite energy. The resulting metric  $\langle \phi \rangle$  is then independent of the choice of reference metric  $\phi_0$ . More generally, in the relative case of a flat morphism  $\mathcal{X} \to B$  between integral schemes of relative dimension n

and an Hermitian line bundle  $L \to \mathcal{X}$  this construction produces an Hermitian line bundle  $\langle L \rangle \to B$  over the base B.

In particular, taking X to be an n-dimensional stable variety  $L := K_X$  one obtains a canonical metric on the complex line  $\langle K_X \rangle$ , induced by the finite energy metric on  $K_X$  determined by the volume form of the Kähler-Einstein metric on the regular locus of X.

Let now  $\mathcal{M}$  denote the moduli space of all n-dimensional stable varieties with a fixed Hilbert polynomial [Kol10, Kol]. Using the existence of a universal stable family  $\mathcal{X}$  (in the sense of Kollar) over a finite cover of each irreducible component of the moduli space one obtains a  $\mathbb{Q}$ - line bundle  $\mathcal{L}$  over  $\mathcal{M}$ , induced by the fiber-wise top Deligne pairings  $\langle K_X \rangle$ . We conjecture that the metric on  $\mathcal{L}$  induced by the fiber-wise Kähler-Einstein metrics is continuous (as in the case of stable curves [FiM12]). Confirming this conjecture would require a more detailed analysis of the dependence of the Kähler-Einstein metric on the complex structure that we leave for the future.

As is well-known the curvature of the corresponding metric over  $\mathcal{M}_0$  coincides (up to a numerical factor) with the Weil-Peterson metric  $\Omega_{WP}$ . In particular, it is strictly positive as a form (in the orbifold sense). Under the validity of the previous conjecture one thus obtains a canonical extension of the induced Weil-Peterson metric on  $\mathcal{M}_0$  to its compactification in  $\mathcal{M}$  as a positive current with continuous potentials. It would also be very interesting to know under which assumptions the extension is *strictly* positive in a suitable sense, for example if it is, locally, the restriction of a Kähler metric? These problems are (e.g. by Grauert's generalization of Kodaira's embedding theorem to singular varities) closely related to the problem of showing that the corresponding line bundle  $\mathcal{L}$  over the moduli space  $\mathcal{M}$  is ample (on each irreducible component) and it should be compared with the recent work of Schumacher [Sch12], where an analytic proof of the quasi-projectivity of  $\mathcal{M}_0$  is given. As shown by Schumacher the Weil-Peterson metric  $\Omega_{WP}$  on  $\mathcal{M}_0$  admits a (non-canonical) extension as a positive current with analytic singularities to Artin's Moishezon compactifaction of  $\mathcal{M}_0$ . But the conjecture above is closely related to the problem of obtaining a *canonical* extension of  $\Omega_{WP}$  to  $\mathcal{M}$  as a positive current with continuous potentials.

Finally, it should be pointed out that the top Deligne pairing used above essentially coincides with Tian's CM-line bundle in this setting (by the Knudson-Mumford expansion and Zhang's isomorphism realizing the Chow divisor as a top Deligne pairing). The ampleness of the induced CM-line bundle over general moduli spaces of K-stable polarized varieties was recently speculated on by Odaka [Oda13].

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